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Contents

Sr. No	Article / Authors Name	Pg No
01	A Visible Light-Responsive Mixed-Valence Bimetallic Eu–Zr MOFBased	1 - 19
	Nanoarchitecture toward Efficacious H2O2 and H2 Production	
	- Srabani Dash,# Suraj Prakash Tripathy,# Satyabrata Subudhi, and	
	Kulamani Parida*	
02	Fault Diagnosis in Chemical Reactors with Data-Driven Methods	20 - 43
	-Pu Du, Nabil M. Abdel Jabbar, Benjamin A. Wilhite, and Costas	
	Kravaris	
03	Celebrating the Birth Centenary of Quantum Mechanics: A Historical	44 - 66
	Perspective	
	-Venkat Venkatasubramanian*	

A Visible Light-Responsive Mixed-Valence Bimetallic Eu–Zr MOFBased Nanoarchitecture toward Efficacious H2O2 and H2 Production

Srabani Dash,# Suraj Prakash Tripathy,# Satyabrata Subudhi, and Kulamani Parida

ABSTRACT

A mixed-valence bimetallic Eu/Zr MOF has been fabricated via a one-step solvothermal method by incorporating Eu3+ ions into the Zr-MOF, thereby making a single-component photocatalyst that can be utilized towardrobust photon utilization from the visible light spectrum for the photocatalytic production of green energy like H2 and H2O2. The one-step synthesized bimetallicEu/Zr-MOF exhibits more visible light captivation properties along withimproved charge carrier separation, confined band gap, and excellent ligand-tometal charge transfer (LMCT) because of the existence of an interconvertibleEu3+/Eu2+ ion pair compared with the pristine MOF counterparts. The addition ofEu ions directed to an upsurge in the electron density around Zr4+ ion, as seenfrom XPS analysis. Moreover, the introduction of Eu3+ enhanced the excitonsegregation, as seen from PL and EIS analyses, thereby leading to superiorcatalytic performances. An increased photocatalytic H2 generation efficacy of 331.26 μ mol h-1 (ACE = 2.42%) was demonstrated by the synthesized EZUNH-2 MOF, which is approximately three times greaterthan pristine MOFs. As a result, the bimetallic EZUNH-2 MOF can be easily utilized as a robust photocatalyst that has increased inclinations to produce H2O2 at 35.2 μ mol h-1, around 4 times more than that of the parent material. Consequently, the one-potsynthesized bimetallic MOF paves a suitable mechanistic pathway for paramount performance toward photocatalytic H2O2 and H2production.

1. INTRODUCTION

The usage of hydrogen peroxide (H2O2) as an environmentally friendly oxidant in industries like chemical synthesis, food andpaper manufacturing, medical decontamination, and wastewater purification has recently attracted a lot of attention from researchers on the front lines.1,2 The water solubility andreleasing H2O as a byproduct are the major factors to makeH2O2 as an energy carrier for future generations.3-5 Severaltechniques have been invented to produce H2O2, including theanthraquinone process, alcohol oxidation, direct synthesis from the mixture of oxygen and hydrogen gases, and electrochemical synthesis. But some of the drawbacks of the synthesis methods listed above include the need for a lot of energy and solvents, as well as the increased risk of explosion brought on by the combination of H2 and O2 gases. An effective and low-energy green method for producing H2O2 is therefore highlydesired.6,7 In recent years, the photocatalytic H2O2 production has received massive attention because of the use of photocatalysts to accomplish the reaction with O2-saturated H2O, alcohol, and light energy, 8-10 In addition, water splitting via a photon-assisted hydrogen evolution reaction has beenperformed to assess the semiconducting materials' photocatalytic energy production capabilities. Following theinnovation of Fujishima and Honda in hydrogen production from H2O utilizing TiO2 semiconductor under light irradiation, which aimed to lower the world's energy demand, the production of H2 from photocatalytic water splitting reaction has developed into an active study area. Moreover, toovercome the current energy crisis, hydrogen (H2) has drawn a lot of interest as a clean, sustainable energy substitute for nonrenewable energy sources. As a result, owing to its lowcost and simple procedure, the photocatalytic hydrogengeneration via water splitting has been considered a promising method. 11–13 Some bimetallic MOFs like Ce–Co MOF, Ni–Ti MOF, and Ce–Zr MOF show excellent hydrogen evolution efficiency as photocatalysts. 14–16

The current situation of an expanding world population, accompanied by the quick industrial development and impending fossil fuel depletion, has prompted efforts to find environmentally favorable and affordable substitute forrenewable energy as well as sustainable catalysis. In order toovercome such issues, the development of a suitable photocatalyst with a tunable band structure, excellent photostability, and greater exciton separation ability is a major challenge.

Zr⁴⁺
BDC-NH,
DMF

Linker
(BDC-NH,)
DMF

Stirring (1 h)

Linker
(BDC-NH,)
DMF

Stirring (1 h)

Linker
(BDC-NH,)
DMF

Stirring (1 h)

Stirring (1 h)

Linker
(BDC-NH,)
DMF

Stirring (1 h)

Stirring (1 h)

Scheme 1. Facile One-Step Synthesis Method of (Eu/Zr) Bimetallic MOF

Since many years, different scientific groups have been enduring multiple semiconducting photocatalysts to get oversuch problems.17-19 Yet, in recent times, organic-inorganichybrids, also known as metal-organic frameworks (MOFs), have become extremely well-liked because of their exceptionalbenefits, which include high specific surface area, flexible functionalization, high porosity, and tailored compositions linked to an infinite variety of organic linkers and metalclusters, resulting in a broad range of potential MOF nanostructures. However, pristine MOFs have some limitations like faster recombination of charge carriers, lack of asuitable band gap, and low light absorption tendency, whichmakes them inadequate photocatalysts. Hence, to conquerthese, some modifications have been involved, such as functional group introduction, guest molecule introduction, and making heterojunction with other photocatalysts.17,18,20However, composite materials have some limitations, likereduction of surface area, undesirable active site coverage, deficient interaction among photocatalysts, lower chargecarrier separation, etc. In recent years, scientists have been observing paths to surmount these difficulties in making composite photocatalysts, which have created an extensive need for the production of single-component MOF-basedphotocatalysts.21,22 Nowadays, single-component MOFbasedphotocatalysts are in high demand, so researchers are exploring avariety of strategies to produce these type of photocatalysts, with tailored functionalities and tuned band structures for awide range of applications.12,23–25Generally, to overcome the above-mentioned issues, various mixed-metal MOFs such as UiO-66-NH2 (Zr/Hf), UiO-66-NH2 (Pt/Sn), UiO-66 (Ti/Zr), MOF-NH2 (Fe/Ti), UiO-66-NH2 (Ce/Zr), etc. were reported.15,26-28 Moreover, thelanthanide-based metal-organic frameworks (Ln-MOFs), especially Eu3+, are interesting due to their adaptable coordination geometry and distinctive luminescent andoptical-electrical properties. In current years, the application of Ln-MOFs as photocatalysts has gained significant interestdue to their unique physicochemical as well as surfacefunctionalization properties. 29,30 They also show easily interconvertible oxidation states of the Europium ion (Eu3+/Eu2+), so the Eu ion insertion presents as a superior alternative over other metals for MOF fabrication. Yet the pristine EuMOF is not stable enough or has insufficient excitonsegregation capabilities; therefore, to achieve optimal photocatalytic activity and stability, it is crucial to optimally introduce these redox mixed valence Eu3+/Eu2+ ions intowater-stable frameworks, such as visible-light-responsive functionalized UiO-66 series Zr-based MOFs. The presence of easily interconvertible Eu2+ and Eu3+ oxidation states in anaqueous stable framework such as visible lightactive UiO-66series Zr-based MOFs promotes superior photocatalyticactivity and stability.32 Since the incorporated metal ionsdetermine their properties, bimetallic MOFs are expected tohave new functionalities in addition to their structural complexity. A second hetero metal node can be added to thesame framework to create synergistic effects that improve itsinherent qualities. 32,33 The integration of Eu3+ into the UiO66-NH2 framework not only builds a stable and porousstructure but also regulates the energy levels of the MOFs. Wehave chosen the Eu3+ ion as the constituent metal due to itslow reduction potential (Ered (Eu3+/Eu2+) = -0.35 V vs NHE) and the ability of the resulting Eu2+ ion to be reconverted easilyto its initial state.34 Furthermore, it is very difficult to find aphotocatalyst that meets the thermodynamic requirements forphotocatalytic water oxidation (At pH = 7, $O2/\bullet O2 - /H2O2 = -0.33$ V and +0.69 V, OH-/OH• = 1.99 V). The key findings for a one-pot synthesized bimetallic Eu/Zr MOFbasednanoarchitecture toward visible light-supported H2O2 production and H2 evolution have been presented in the currentinvestigation. After extensive research on the neat Zr-MOF and its bimetallic MOF, such as (Eu/Zr) UiO-66-NH2 [EZUNH], the composite showed improved exciton segregation androbust LMCT, accompanied by the mixed valency states of Eu3+ and Eu2+, which endorse its excellent photocatalyticoutput. Moreover, due to the strong bond of Zr-O, the pristing MOF with Zr as the metal center has high stability toward acid—base, aqueous, and thermal changes, and a similar impact can be inherited in the EZUNH MOF.35 The currentwork represents the one-pot synthesized EZUNH bimetallicMOF-based nanomaterials toward photocatalytic H2O2 and H2production. Among the prepared bimetallic MOFs, the EZUNH-2 exhibits boosted photocatalytic H2 production upto 331.26 μmol h-1 and H2O2 production up to 35.2 μmol h-1, which is around 3-fold and 4-fold higher than that of the pristine MOF.

2. MATERIALS AND METHODS

2.1. Chemicals Used. Several chemicals, for example, zirconium chloride (ZrCl4, 99.99%), 2-aminoterephthalic acid(BDC-NH2, 99%), europium(III) chloride hexahydrate(EuCl3.6H2O, 99.99%) Nafion-117 (~5% in lower aliphaticalcohol/water mixture), and potassium bromide (KBr, 99.5%) were bought from Sigma-Aldrich. In addition, sodium sulfate(Na2SO4, 99%), methanol (MeOH, 99%), isopropanol (99%),

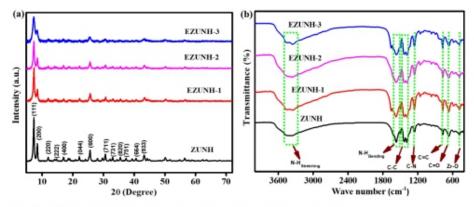


Figure 1. (a) XRD pattern and (b) FTIR spectra of ZUNH and EZUNH-1, 2, and 3 MOFs.

and N,N-dimethylformamide (DMF, 99.8%) were purchased from Merck. All the chemicals involved were used in several reactions.

- **2.2. Fabrication of Photocatalyst.** 2.2.1. Pristine MOFs. The UiO-66-NH2 MOF was prepared by following asolvothermal approach by procuring an equimolar ratio ofmetal salts (ZrCl4, 2 mmol) and linkers (ATA, 2 mmol), asdescribed earlier. 36,37 First, with DMF (40 mL), the metal saltand linker were added individually with continued stirring for 1h. Then, both solutions were mixed together and further stirredfor 1 h. Afterward, the mixture solution was transferred into a Teflon-lined stainless steel vessel and kept for solvothermaltreatment (120 °C for 24 h). After the reaction time was over, the reactor vessel was allowed to cool down to roomtemperature. Subsequently, the pore activation of the productwas done by a solvent exchange method using methanol for 24h. The products were then collected through centrifugation and dried overnight at 80 °C. Last, the obtained yellow-colored sample was named ZUNH. 2.2.2. Bimetallic Eu/Zr-UiO-66-NH2 MOF. The Eu/Zr bimetallic MOF was fabricated in an analogous way to the parent ZUNH MOF through a one-step solvothermal method (120 °C for 24 h), as illustrated in Scheme 1. Herein, this process differs from the ZUNH fabrication method by using mixed metallic salts with variable molar concentrations as Eu0.2 mmol/Zr-1.8 mmol, Eu-0.4 mmol/Zr-1.6 mmol, and Eu0.6 mmol/Zr-1.4 mmol, which were termed as EZUNH-1, EZUNH-2, and EZUNH-3, respectively. Then, to each ofthese salt mixtures, DMF (40 mL) was added. Further, themetal salts and linker were added individually with continuous stirring for 1 h. Then, after both solutions were mixed together, they were further subjected to stirring for 1 h. Afterward, themixture solution was transferred into a Teflon-lined stainlesssteel vessel and kept for solvothermal treatment (120 °C for 24h). After the reaction time was over, the reactor vessel was allowed to cool down to room temperature. Subsequently, the pore activation of the product was done by a solvent exchangemethod using methanol for 24 h. The detailed characterizationmethods involved and other experimental processes followed in this work are reported in the Supporting Information(Experimental Techniques).
- 2.2.3. Photocatalytic H2O2 and H2 Production. The synthesized samples were subjected to analysis of the photocatalytic activity toward H2O2 production under an O2-saturated atmosphere with 2 h of visible light illumination ($\lambda \ge 420$ nm). A suspension solution was prepared by adding 19 mL of deionized water (DI) and 1 mL of isopropanol (IPA) with 20 mg of photocatalyst, and then, the suspension underwent anultrasonication process for about 10 min for proper dispersion of the contents. Subsequently, the solution was kept under O2purging for 30 min in the presence of light to attain an O2saturated atmosphere. After the reaction time, a clear solutionwas obtained by centrifugation of the suspension solution. Thereafter, to 1 mL of the resulting solution, 2 mL of 0.1 M KIsolution and 0.05 mL of 0.01 M ammonium molybdatesolution were added to change the colorless sample into alight-yellow color. Finally, the concentration of the producedphotocatalytic H2O2 was evaluated through a Uv-visible spectrophotometer

Furthermore, the prepared nanomaterials were toward the photocatalytic evolution of hydrogen gas. In this process, a closed quartz batch-type reactor (100 mL) was used to acquire the photocatalytic H2 production efficacy of thefabricated photocatalysts, such as pure ZUNH and bimetallicEZUNH MOFs. Here, 20 mg of the as-synthesized photocatalysts was taken in the photoreactor with 20 mL of 10% V/V MeOH-water mixture, and the visible light source (Xenonarc lamp, 300 W, $\lambda \ge 420$ nm) was irradiated for 1 h. The substances in the reactor were continuously stirred to promote a uniform distribution and avoid particle aggregation during the reaction time. Formerly, by utilizing the Xe lamp, suspension mixture was thoroughly bubbled for 30 min under N2 gas to eradicate the dissolved gases present. The gaseous mixtures that emerged were agitated via direct water displacement and investigated using gas chromatography/GC (GC-7890B, Agilent Technologies) tailored with 5 Åmolecular sieves and a thermal conductivity detector (TCD). Photocatalytic experiments were performed in triplicate tominimize experimental errors.

3. RESULTS AND DISCUSSION

3.1. Physicochemical Characterizations. Figure 1a illustrates a PXRD (powder X-ray diffraction) pattern analysis, which was executed to construe the formation and crystallographic nature of all the synthesized photocatalysts. Typically, the diffraction pattern of the pristine MOF ZUNH suggestshigh crystallinity. Remarkably, the EZUNH MOFs show thecharacteristic XRD pattern, suggesting the conservancy of crystallographic peaks of the pristine grid framework upon theinsertion of Eu ions. The Zr4+ and Eu3+ competitively coordinate with the ligand to produce a disrupted and slightly reduced crystallinity in bimetallic EZUNH MOFs.

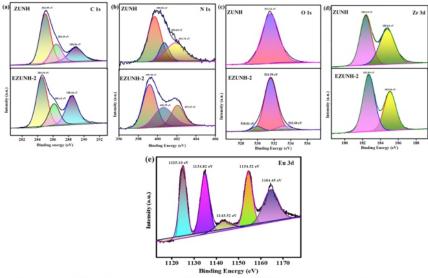


Figure 2. XPS spectra of (a) C 1s, (b) N 1s, (c) O 1s, (d) Zr 3d, and (e) Eu 3d of ZUNH and bimetallic EZUNH-2 MOF.

Consequently, the diffraction at 7.35° and 8.41°, corresponding to the (111) and (200) planes, respectively, gradually decreases as the amount of Eu3+ ions increases. Additionally,the diffraction peaks shown in Figure 1a exhibit a reduction in the intensity of peak, accompanied by broadening of peaks with respect to the parent ZUNH, especially for the two noticeable Zr-cluster peaks, signifying the Eu ions' insertioninto the lattice of MOF and eventually instigating the slightreduction in Zr-Oxy clusters.15,38Additionally, the FT-IR (Fourier transform infrared) studywas performed to detect various functional groups by means oftheir modes of vibration existing in the pristine ZUNH and bimetallic EZUNH MOF, as illustrated in Figure 1b. The observed peaks at 3475 and 3340 cm-1 represent the -Nh2groups of both ZUNH and EZUNH, which are accompanied by asymmetrical and symmetrical vibrational stretching modes, respectively. In the frequency region, the 1640 and 1245 cm-1bands determine the vibrational bending of N-H and stretching of the C-N bond of the aromatic -NH2 groups existing in the ATA linkers, respectively. Furthermore, the small vibrational band at 1511 cm−1 represents the C�Cmoiety of the benzene ring, whereas, the vibrational peaks observed around 485, 763, 662 cm−1 are accompanied by metal-(OC) asymmetric stretching, C♦C stretching of thearomatic ring, and car boxylate O&C&O bending, respectively.12,39 Typically, from the FT-IR analysis results of bimetallic EZUNH-1, 2, and 3 MOFs, it has been observed that the nature is quite similar to pristine ZUNH. Hence, theanalysis outcome suggests an equivalent chemical bondingatmosphere and indistinguishable functional groups in the frameworks along with it signify no

noticeable difference in the peak positions. The FT-IR results support the earlier reported results.32,40 Textural properties of the synthesized MOFs were persistent through the (BET) surface area technique with the support of the N2 adsorption/desorption isotherm method. Figure S1depicts the bimetallic EZUNH-2 MOF's BET adsorption/desorption curve, which is similar to the Type-IV isotherm thatindicates the microporous and mesoporous nature of theframework. Herein, the insertion of Eu ions into the framework ZUNH assisted in the diminution of the BET surface area,i.e., 533.5 m2 g-1 as given in Figure S1. However, thephotocatalytic yield does not distress with the reduction insurface area due to the synergistic effect of bimetallic ions.Regardless, the insertion of Eu ions does not affect the rigidity of the framework structure, which shows that the isothermpattern of the bimetallic MOF persisted similarly to that of theparent ZUNH MOF.18

Chemical states and surface elemental compositions of the prepared photocatalysts were investigated through X-rayphotoelectron spectroscopy (XPS). Herein, for EZUNH-2bimetallic MOF, the XPS survey spectra have been depicted in Figure S2, which validates the presence of Eu, Zr, C, N, and O, as confirmed by the EDAX and elemental mapping analysis. Moreover, from the XPS analysis of the fabricated bimetallic EZUNH-2 MOF, the presence of each element's spectra was compiled and deconvoluted, as illustrated in Figure 2b–f. From the deconvoluted spectra, the C 1s present in EZUNH-2exhibits peaks at 284.54, 285.12, and 288.44 eV, which resemble the carbon atoms of the BDC-NH2 linker that correspond to C&C, C–NH2, and O&C–O, respectively. 13 Also, for N 1s, the deconvoluted XPS peaks observed at

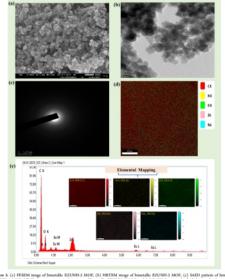
	Binding Energy in eV					
Element					_	Ref.
Carbon/C 1s						
ZUNH	284.99		286.30	288.96		41,42
20% EZUNH	284.54		286.12	288.44		
Speculation:	C=C of linker		C-NH2 of linker	(O=C-C	O) of linker	
Difference:	-0.45		-0.18	-0.52		
			Nitrogen/N 1s			
ZUNH	399.52		400.63	401.76		13
20% EZUNH	399.18		400.59	402.03		
Speculation:	-NH2 of linker		-NH3+ of linker	=NH2+ of linker		
Difference:	-0.34		-0.04	+0.27		
			Oxygen/O 1s			
ZUNH			532.14			43
20% EZUNH	530.01		531.59	533.30		
Speculation:	Mn-O bond		Zr-O bond	Adsorbed H ₂ O		
Difference:	+0.69		-0.55			
			Zirconium/Zr 3d			
ZUNH	183.13		185.61			18,44
20% EZUNH	182.69		185.04			
Speculation:	Zr4+(3d5/2)		Zr4+(3d3/2)			
Difference:	-0.44		-0.57			
			Europium/Eu 3d			
ZUNH						23,45,4
20% EZUNH	1125.10	1134.82	1143.52	1154.52	1164.45	
speculation:	Eu (3d _{5/2})	Sat. Peak		Eu (3d _{3/2})		
Difference:						

399.18, 400.15, and 402.03 eV relate to the -NH2 functionalized groups of linkers such as -NH2 and -NH3 +,respectively, which is almost analogous to the previouslyanalyzed N 1s spectra for pristine ZUNH MOF. Besides, thedeconvoluted O 1s XPS spectra found at 530.01, 531.59, and533.30 eV signify lattice O, metal-O, and surface-adsorbedH2O molecule, respectively, for the prepared bimetallicEZUNH-2 MOF. Additionally, the deconvoluted peaks forZr found at 182.69 and 185.04 eV represent the two spin states of Zr, such as 3d5/2 and Zr 3d3/2 in the EZUNH-2 sample. Thedownshifting of Zr spin states binding energy demonstrates that the electron density of Zr4+ increases due to the transfer ofelectrons from lower oxidation species (Eu3+) in the bimetallicEZUNH-2 MOF. In the deconvoluted peaks, the XPS of the Eu peak comprises dual sets corresponding to the states Eu 3d5/2

and Eu 3d3/2. Herein, the deconvoluted peaks of Eu correspond to the 3+ oxidation state of Eu 3d5/2 and Eu 3d3/2, respectively; also, the presence of the 2+ oxidation state of Euions has been confirmed from the plotted data. From the deconvoluted plot of Eu, the existence of both the 3+ and 2+oxidation states in the MOF has been clearly perceived, suggesting the formation of mixed valency MOF. Moreover, the XPS analysis outcome confirms the formation of a singlecomponent bimetallic EZUNH MOF along with it unveils therole of Eu3+/Eu2+ redox pair. The deconvoluted peak values of XPS confirm the pristine ZUNH and bimetallic EZUNH-2, along with their corresponding chemical environments, are presented in Table 1 for lucid comprehension. Additionally, the ICP-OES analysis result confirms the existence of Eu:Zr in the bimetallic MOFs with different molar ratios, such as EZUNH-1 (0.08:0.84), EZUNH-2 (0.17:0.76), and EZUNH-3(0.26:0.67), respectively

3.2. Morphological Studies. To study the surface morphology and elemental composition of the pristineZUNH and one-pot synthesized mixed-metallic EZUNH-2MOF, the FESEM (field-emission scanning electron microscopy) and HRTEM (high-resolution transmission electronmicroscopy) analyses were carried out. The FESEM image(Figure 3a) reveals the morphology of EZUNH-2, which isquite similar to that of the pristine MOF, as shown in FigureS3a. Additionally, the analogous morphological features of themixed-metallic MOF substantially corroborate the PXRDresults, indicating a framework structure similar to that ofpristine MOFs. From the HRTEM images given in Figure 3b, the octahedral morphology of EZUNH-2 was distinctly perceived at 20 nm scales as depicted. Also, it is significantly observed that the morphology of the bimetallic MOF is well supported by the pristine MOF (Figure S3b). Herein, the addition of Eu into the Zr-framework significantly does not affect the crystal structure, which specifies the similarmorphology of the synthesized bimetallic MOF. The parentZUNH MOF has high sensitivity toward the electron beams; also, the bimetallic framework exhibits a similar property. However, it is difficult to find high-quality HRTEM images aswell as a welldefined SAED pattern. The SAED analysispattern given in Figure 3c corresponds with the neat ZUNHMOF, as reported formerly by us.17 Furthermore, the elemental color mapping (Figure 3e-j) and EDS (Figure 3d) outcomes offer extra sustenance to the existence of C, O, N, Eu, and Zrelements in the EZUNH-2 bimetallic framework, which is also confirmed by XPS analysis.

3.3. Optical Characterization. The optical properties of the prepared photocatalysts were analyzed through Uv-visiblediffuse reflectance spectra (UV-DRS) and are shown in Figure 4a. The two intense bands obtained at 265 and 365 nm were



associated with ZUNH, which are attributed to the lone pair electron n- π^* transitions present in the -NH2 group of the ATA linker and the $\pi^-\pi^*$ transitions overlapping the ATA linker accompanied by Zr–O cluster absorption bands,respectively.15,47 Additionally, the similar bands in the bimetallic EZUNH MOFs imply the suitable introduction of Eu ions into the ZUNH framework structure. The optical band gaps of the materials were evaluated by following the Kubelka–Munk equation, as given below in eq 1.

$$\alpha h \nu = A(h\nu - E_g)^{n/2}$$
(1)

Herein, α indicates proportionality constant, h represents Planck's constant, v specifies the frequency of incident light, Asignifies the light absorption coefficient, and Eg represents the

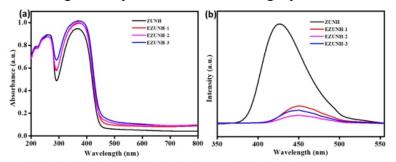


Figure 4. (a) UV-visible DRS spectra and (b) PL spectra of pristine ZUNH and bimetallic EZUNH MOFs.

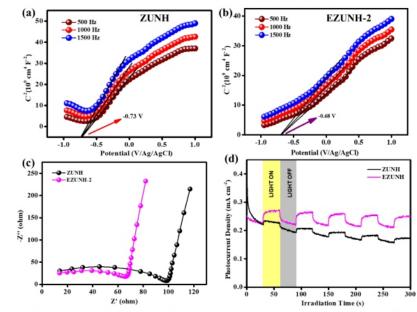


Figure 5. (a-b) Mott-Schottky plot, (c) EIS analysis plot, and (d) transient photocurrent representation of ZUNH and EZUNH-2 MOFs.

optical band gap, respectively. Additionally, the parameter "n"describes the probable reduction of electronic transitions withn = 1, 3, 4, and 6, which correspond to allowed directtransition, forbidden direct transition, allowed indirecttransition, and forbidden indirect transition, respectively. Herein, the fabricated mixed metallic MOF and pristineMOF exhibited allowed indirect transition (n = 4).16 Theoptical band gap of ZUNH was calculated as 2.67 eV. Also, the calculated band gap energy for the synthesized EZUNHbimetallic MOFs were 2.63, 2.59, and 2.52 eV for EZUNH-1, EZUNH-2, and EZUNH-3 by varying Eu percentage, respectively. The Tauc plots of the prepared photocatalysts are given in Figure S4. It is fairly evident that the addition of Eu ions in their mixed valency states (Eu2+ and Eu3+) accelerates the process of excellent ligand-to-metal charge transfer (LMCT) and that the enhancement of their light absorption tendency is a result of the reduction of the optical band gap, which

in turn improves photocatalytic performance.38

The photoluminescence (PL) analysis was carried out to reveal the migration and separation efficacy of photogenerated excitons. Usually, the lower intensity of PL peaks indicates superior separation of exciton pairs, which results in higherphotocatalytic activity. From the resulting PL spectra plot, EZUNH-2 shows substantially weaker fluorescence emission intensity compared to the pristine ZUNH and the other metal-doped composites such as EZUNH-1 and EZUNH-3, which explains the lowering in the recombination rate of photogenerated excitons, as illustrated in Figure 4b. Also, the Zr4+ has been substituted by some Eu3+ content, which exhibits a superior electron trapping tendency with a slight red shift of the emission peak observed that is associated with then arrowing effect of Eu doping. 48,49 Moreover, the time resolved photoluminescence (TRPL) study of the neat ZUNH and bimetallic EZUNH-2 MOF was used to investigate the exciton lifetime. This analysis is illustrated in Figure S5 and Table S1, which are adapted to a biexponential model equation, as stated in eq 2.

$$R(t) = A_1 \exp\{-t/\tau_1\} + A_2 \exp\{-t/\tau_2\}$$
(2)

where A stands for amplitude, τ is the exciton life span of each individual component, and R is the normalized emissionintensity. Here, the extended lifetime $\tau 2$ and short lifespan $\tau 1$ reflect the nonradiative relaxation mechanism of light excitons and radiative recombination, respectively. Herein, the averagelifetime (τavg) can be evaluated to explicate the whole TRPL characteristics of both exponential decays by using eq 3.

$$\tau_{\text{avg}} = \frac{A_1 \tau_1^2 + A_2 \tau_2^2}{A_1 \tau_1 + A_2 \tau_2}$$
(3)

For ZUNH and EZUNH-2 MOFs, the excited state lifetimes have averages of 0.504 and 0.533 ns, respectively. This resultincreases the photocatalytic efficiency and supports the PLexperimental results by explaining EZUNH-2 longerexciton than due charge antirecombination.37,50

Electrochemical The Mott–Schottky (MS) study was performed to investigate the band position and charge transfer path in the mixed metallic MOF, as illustrated in Figure 5a,b. The obtained flat band potential (Efb) from the MS plot is a fundamental parameter to obtain the band edge potentials and paths followed by the photogenerated e-/h+ in a semiconductor photocatalyst. Here, the flat band potentials of ZUNH and EZUNH-2 are-0.73 and -0.68 eV, respectively, in contrast to the Ag/AgClelectrode, which was evaluated by extrapolating the C2- = 0 curve. Subsequently, the CB of ZUNH and EZUNH-2 were found by following eq 4.

$$E_{\text{(NHE,pH=7)}} = E_{\text{Ag/AgCl}} - 0.059 \times (7\text{-pH of the electrolyte}) + 0.198$$
(4)

In MS analysis, the pristine ZUNH and bimetallic EZUNH2 show a positive slope, which endorses the n-type behavior of both the prepared MOFs. Furthermore, it is relatively evidentthat the n-type characteristic of the neat ZUNH was inherited in the bimetallic MOF. Moreover, the slope of the MS graphwas inversely proportional to the charge carrier density. So, thereduction in the slope of EZUNH-2 was supported by higher carrier density, which was calculated using the eq 5.

$$\frac{1}{C^2} = \frac{2}{q \in \in_0 N_d} \left(E - E_{fb} - \frac{K_b T}{q} \right) (\text{n-type})$$
(5)

where kb represents the Boltzmann's constant, E signifies the applied potential, T is the absolute temperature, C shows the space charge capacitance, Nd is the donor density, Na exhibits acceptor density, and q represents the electronic charge. $\epsilon 0$ is the permittivity in vacuum, and ϵ is the dielectric constant of the photocatalyst. A positive shift of 0.05 eV was observed in the flat band potential of EZUNH-2 with respect to pristine ZUNH, signifying the enhanced charge carrier density in the mixed metallic MOF (containing the Eu redox couple). Hence, this higher carrier density promotes the boosted photocatalytic result in the bimetallic EZUNH-2 MOF. Moreover, from the Efb values, the band structure of parent ZUNH (VB = 2.03 eV,CB = -0.64 eV) and bimetallic EZUNH-2 (VB = 2.0 eV, CB = -0.59 eV) vs Ag/AgCl was obtained by using eq 5 in the NHEscale. Furthermore, the VB and CB values were changed in the NHE scale and are illustrated in mechanism Scheme 2. It has been evidently observed that the CB is slightly shifted toward VB in EZUNH-2 before a minor variation in the actual VB position, pointing toward a robust and effective LMCT in EZUNH-2 MOF rather than pure ZUNH.51

Also, from electrochemical impedance spectroscopy (EIS)analysis, the effective separation and migration of chargecarriers in the mixed metallic MOF were discovered. Basically,the smaller lower interfacial charge transfer resistance refers to the semicircular arc diameter, which exhibited superior transfer of charge carriers in the material. Figure 5c exhibits the EZUNH-2 composite that shows a quite reduced semicircular arc diameter compared to the pristine ZUNH MOF that recommends the antirecombination of charge carriers, which indicates improved electrical conductivity of the bimetallic EZUNH-2 MOF. 50,52,53 The EIS outcome of the bimetallic EZUNH-2 MOF was well supported by the PL studies. Inaddition, current versus potential measurements (LSV) were performed to analyze the photogenerated charge carrier transfer and the mechanism of the photocatalyst. This investigation was performed for the parent ZUNH and all the bimetallic EZUNH MOFs at 5 mV s-1 in an appropriate potential range, as depicted in Figure S6. The pristine ZUNHMOF produces an anodic photocurrent, which represents an ntype feature. Also, the bimetallic EZUNH MOFs exhibit analogous characteristics, possessing enhanced photocurrent compared to the pristine MOF

At last, the transient photocurrent investigation was executed for the ZUNH and EZUNH-2 MOFs. The experiment was accomplished under alternative cycles in darkand visible light ($\lambda \ge 420$ nm) irradiation environments to exhibit the boosted separation efficiency of exciton pairs. Herein, the analysis result was well established with the production of photocurrent, which mainly comprises the diffusion of photogenerated e–s to the back contact, and the consumption of h+s takes place by the hole scavengers present in the electrolyte solution. As indicated in Figure 5d, the EZUNH-2 MOF exhibited increased transient current, which signifies an enhanced lifetime of the excitons in the bimetallic MOF than in the pristine MOF. This was accredited to the effective separation of photoexcitons or lower recombination rate promoted through the Eu ion insertion into the ZUNH framework, and this eventually improves the photocatalyticactivity. Moreover, the transient photocurrent investigation outcomes validate the PL and EIS results, as reported earlier.

4. PHOTOCATALYTIC PERFORMANCE

The photocatalytic behaviors of the prepared samples were analyzed by executing the hydrogen peroxide (H2O2) andhydrogen (H2) production reactions under visible lightirradiation. Primarily, the H2O2 production reaction was executed in an O2-saturated atmosphere with 2 h of visible light ($\lambda \ge 420$ nm) illumination under ambient conditions. Butthere was no H2O2 production observed in the absence of acatalyst or light, which explains that the catalysts and light are

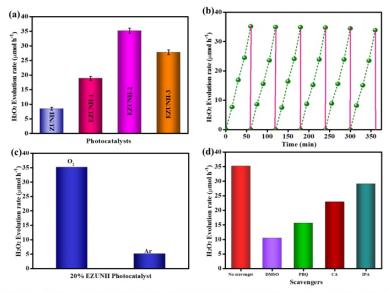


Figure 6. (a) H_2O_2 production rate of pristine ZUNH and bimetallic EZUNH-1, 2, and 3 MOFs (for each photocatalyst, the H_2O_2 production is represented as mean \pm SD), (b) reusability test for H_2O_2 production over EZUNH-2 MOF, (c) H_2O_2 production rate in different environments, and (d) scavenger test for H_2O_2 production rate.

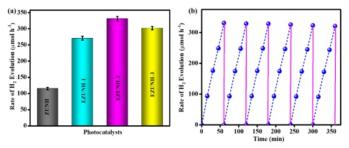
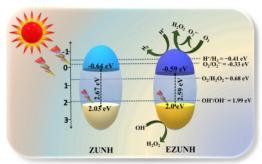


Figure 7. (a) H_2 production rate of pristine ZUNH and bimetallic EZUNH-1, 2, and 3 MOFs (for each photocatalyst, the H_2 production is represented as mean \pm SD) and (b) reusability test for H_2 production over EZUNH-2 bimetallic MOF.

he key factors for the reaction to occur. The H2O2 production rate of pristine ZUNH was found to be $8.52~\mu$ mol h=1. Herein,among the prepared bimetallic photocatalysts, the EZUNH-2MOF shows maximum H2O2 production rate (35.2 μ mol h=1),as shown in Figure 6a. Similarly, a series of experiments wasperformed for H2O2 production of EZUNH-1 and 3 MOFs,where the obtained outcomes were 27.8 and 18.9 μ mol h=1respectively. Among all the photocatalytic H2O2 production results by the prepared materials, it was observed that the bimetallic MOFs show significantly greater photocatalyticactivity than the pristine ZUNH, which was attributed to their enhanced light absorption behavior and charge carrierantire combination capacity due to the presence of Eu ions.

The photocatalytic activity of the synthesized bimetallic MOF (EZUNH-2) was four times greater than that of the pristine ZUNH MOF. Furthermore, the reusability experiment of thematerial indicates that the EZUNH-2 MOF exhibits photostability for up to four successive cycles, as depicted in Figure 6b. Additionally, the O2 dependency was studied for the photocatalytic H2O2 production by EZUNH-2 bimetallic MOF, and under Ar and O2 gas purging, two separatereactions were executed, which resulted in a very small quantity of H2O2 production taking place in the Aratmosphere, as shown in Figure 6c. Thus, it appears that the existence of an O2 atmosphere is necessary for the photocatalytic generation of H2O2. Moreover, the impact of

Scheme 2. Schematic Illustration of the Mechanistic Path for Photocatalytic H₂O₂ and H₂ Production by the Prepared Bimetallic MOF



scavengers affecting the production of H2O2 is shown in Figure 6d and is discussed briefly in the Supporting Information. The comparison table to show the significance of the prepared bimetallic MOF for H2O2 production is given in Table S4. Alternatively, the photocatalytic H2 evolution efficiency of the synthesized materials was also measured with visible lightirradiation. In the absence of light or catalyst, the blankreadings were taken to prove that both the catalyst and lightirradiation are essential for the H2 evolution reaction to takeplace. The pristine MOF ZUNH exhibits less H2 evolution rate of up to 115 µmol h−1, which is due to the fasterrecombination of exciton pairs. However, the bimetallicMOFs show an enhanced H2 production rate because of the superior visible light absorption along with greater chargesegregation and transfer. The photocatalytic rate of H2evolution for the prepared bimetallic MOFs EZUNH-1, 2, and 3 was found to be 331.26, 302.01, and 270.81 µmol h-1, respectively, as depicted in Figure 7a. From the resulting plot, the EZUNH-2 exhibits the utmost photocatalytic H2production rate, which is almost 4-fold greater than the pristing ZUNH MOF. The apparent conversion efficiency(ACE) was analyzed to be 2.42%, as shown in Table S3. Thestability of the bimetallic MOF EZUNH-2 was checked byperforming four consecutive cycles of H2 evolution, which suggests that there was no substantial change in the rate of H2production, as shown in Figure 7b. Also, the postphotocatalyticXRD has been studied, as given in Figure S5. Table S4 suggests the comparative H2 evolution table to show the importance of the prepared EZUNH-2 bimetallic MOF.

5. MECHANISM INSIGHT

Herein, Scheme 2 demonstrates the plausible mechanism of the fabricated mixed metallic EZUNH-2 MOF towardphotocatalytic H2O2 production and H2 evolution undervisible light irradiation, which was expressed thoroughlybased on the aforementioned analytical outcomes. Toinvestigate the fundamental mechanism, characterization analyses like UV-visible DRS, PL, XPS, MS, and EIS weretaken into consideration, from which the experimental results indicated a suitable reduction of band edge as well as superior exciton pair segregation by the insertion of Eu ions into the ZUNH framework toward the H2O2 and H2 production.

Mainly, the reduction in band gap was observed for the EZUNH-2 MOF, which was effectively obtained from Uv-visible DRS and MS investigation outcomes, signifying a robustLMCT and excellent light-harvesting propensity compared to the pristine counterparts. From the Tauc plot, the obtained band gap of bimetallic EZUNH-2 MOF was 2.59 eV, which is lower than the pristine ZUNH.42 The Eu ion insertion into the pristine MOF lattice reduces the band gap and increases the visible light captivation capacity. Also, the analysis outcomes of PL and EIS offer a fascinating observation of improvedantire combination of photoexcitons in the prepared bimetallic MOFs and promote a longer exciton lifespan than the pristine ZUNH. The Eu3+/Eu2+ redox pair shows amended LMCT, which is well supported by the reduction of Zr4+ binding energy in XPS peaks, signifying superior transfer of electrons in the cluster,

accelerating the boosted photocatalytic activities.15,53 When the prepared materials were kept for visible light irradiation, the electrons from VB got excited to CB, leaving behind the holes and leading to the production of hydroxyl radicals. These radicals act as strong oxidants for H2O2 production. This process leads to the undesirable recombination of exciton pairs, which rapidly reduces the photocatalytic activity. As reported by different groups, the Eu3+ ions have unfulfilled 4f orbitals, which lead to thereduction of Eu3+. Herein, the Eu3+ may accept electrons in theCB to form Eu2+, and these Eu2+ ions transport the electrons to dissolve O2 to produce superoxide radicals, inhibiting therecombination of photoexcited charge carriers. This shows that Eu3+ acts as an electron scavenger. However, it has beenobserved that the excessive presence of europium ions may act as recombination centers for photogenerated e- and h+ or elsemay block the active sites of the catalyst's surface. Therefore, increasing the dopant concentration causes a sharp decrease inphotoactivity, as seen for EZUNH-3.31,48,54,55 The MS analysisoutcome determines the CB and VB positions of EZUNHMOF as 2.0 and -0.59 eV, respectively. Upon light irradiation, the bimetallic EZUNH-2 got excited to produce the photoexcitons, which occur by the electrons transferred from VB to the CB, leaving behind holes in the VB, which directed to the probable efficient reduction and recombination because of the existence of mixed valency metal ions in the framework structure. To investigate the photocatalytic mechanistic insights of H2O2 production, the CB potential of EZUNH-2 MOF was found to be -0.59 eV, which is more negative than that of the potential of H2O2 production at -0.33 eV, which is the one-electron pathway, and at 0.68 eV, which represents the two-electron pathway. 56 Thus, the longer lifetime of excitedelectrons endorses O2 reduction to produce H2O2, as shown ineqs 6 and 7.

Single-step two-electron reduction pathway;

$$O_2 + 2H^+ + 2e^- \rightarrow H_2O_2$$
 (6)
Two-step single-electron reduction pathway;

$$O_2 \rightarrow {}^{\bullet}O_2^- \rightarrow H_2O_2$$
 (7)

Although, the VB of EZUNH-2 consisted of photogenerated h+s, which further get trapped by ethanol (EtOH) as asacrificial agent. Also, the VB potential of EZUNH-2 satisfies the formation of hydroxyl radicals (OH \bullet /OH- = 1.99 eV vsNHE). Henceforth, the combination of two OH \bullet radicals produces H2O2, as depicted in eq 8.

$$OH^{\bullet} + OH^{\bullet} \rightarrow H_2O_2$$
 (8)

The efficient H2O2 production shown by EZUNH-2 bimetallic MOF is 35.2 μmol h–1. Moreover, the reusability test was performed for stability check of the as-synthesizedmaterial, which endows the stability of the sample up to foursuccessive cycles, as displayed in Figure 6b. Furthermore, toexplicate the mechanistic pathway, scavenger tests were carriedout. As depicted in Figure 6d, there are different scavengingagents like isopropanol (IPA), dimethyl sulfoxide (DMSO),parabenzoquinone (PBQ), and citric acid (CA) were used toexplore the role of OH•, e–, •O2–, and h+, respectively,throughout the photocatalytic H2O2 production reaction. Theobserved substantial decrease in the H2O2 productionefficiency with PBQ and DMSO addition during the reactiontime can be witnessed in Figure 6d, which demonstrates thate– – and •O2– act as the major active species for O2 reduction. In addition, IPA and CA play a very certain role towardphotocatalytic H2O2 production, which signifies that OH• andh+ are less reactive for the O2 reduction reaction. Therefore, the analysis results of scavenger tests follow the order e– > • O2– > OH• > h+ toward H2O2 production through O2reduction. 9,13 Also, supporting the scavenger tests, the TA(terephthalic acid) and NBT (n itroblue tetrazolium hydrochloride) test results confirm the formation of •O2– and OH• radicals by the bimetallic MOF

EZUNH-2, as shown in Figure S7.57,58

Moreover, the photocatalytic H2 evolution by mixed metal MOFs has been studied by following the mechanistic pathway. With the hole scavenger (10% methanol v/v solution), themanufactured mixed metal composite EZUNH-2 wasemployed for the photoreduction of water to create H2 gas. The H2O molecules can readily be adsorbed on the MOFsurface due to the hydrophilic character of the EZUNH-2 framework, which is enhanced by the presence of –NH2moieties. Following this, the protons formed by the dissociation of water molecules uptake the photoexcitedelectrons from the CB of EZUNH-2 and get reduced togenerate H2 gas. The reactions that followed for photocatalyticH2 production were designated as follows:

EZUNH-2 + h
$$\nu$$
 → EZUNH-2 [VB(h⁺)/CB(e⁻)] (9)
 $H_2O \rightarrow H^+ + OH^-$ (10)
EZUNH-2 [CB(e⁻)] + $H^+ \rightarrow H_2$ (11)
 $CH_3OH + h^+ \rightarrow {}^{\bullet}CH_2OH + H^+$ (12)
 ${}^{\bullet}CH_2OH \rightarrow CH_2O + e^- + H^+$ (13)
 $2H^+ + e^- \rightarrow H_2$ (14)

6. SUMMARY

Summarily, the synthesized single-component mixed metallic EZUNH MOF has been found to be a robust photocatalysttoward H2O2 and H2 production under light irradiation. The superiority of the current work is followed by (i) The prepared single-component photocatalyst is a mixed valency (Eu3+/Eu2+) and bimetallic (Eu/Zr) MOF, i.e., (Eu3+/Eu2+) EZUNH, which is synthesized by a onestep solvothermal method.(ii)The integration of mixed valency (Eu3+/Eu2+) into the bimetallic framework gives superior exciton pair antirecombination, high light absorption efficacy, and enhanced electrochemical properties, resulting in boosted photocatalytic activity compared to the pristine ZUNH MOF. (iii) The joint effect of both the metal nodes and -NH2 functionalized linker in EZUNH MOFs results in a blueshift of the band gap and effective band edge positions for superior light absorption with a slight reduction in surface area. (iv)The fabricated EZUNH-2bimetallic MOF exhibits the highest photocatalytic H2evolution rate of up to 331.26, along with an ACE value of 2.42%, which is almost 3fold that of the pristine MOF. Also, the photocatalytic H2O2 production rate has been found to be35.2, which is up to 4-fold compared to that of the pristine framework. Herein, e-s are the main active species for H2O2 production, as found from the scavenger test experiments. Finally, from the above results, it has been confirmed that the bimetallic EZUNH MOF is the first used photocatalyst toward H2 and H2O2 production (μ mol h-1) under visible lightirradiation.

■ ASSOCIATED CONTENT

*si Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.iecr.4c04234.1. Experimental techniques (1.1. CharacterizationTechniques, 1.2. FTO preparation, 1.3. Scavenger testprocedure, and 1.4. TA and NBT analysis procedure);BET surface of EZUNH-2 (Figure S1); XPS surveyspectrum of EZUNH-2 (Figure S2); (a) FESEM and(b) HRTEM images of pristine ZUNH MOF (FigureS3); Tauc plot for ZUNH, EZUNH-1, EZUNH-2 and EZUNH-3 (Figure S4); TRPL plot of EZUNH-2 (Figure S5); LSV plot for pristine ZUNH and bimetallicEZUNH-1,2 and 3 MOFs (Figure S6); TA and NBTanalysis results (Figure S7); postphotocatalytic and 24-haqueous string PXRD (Figure S8); TRPL data of ZUNHand

EZUNH-2 (Table S1); apparent conversionefficiency (ACE) expression and calculation for photocatalytic hydrogen evolution by the prepared photocatalysts (Table S2); comparative table for photocatalytic H2 production by various photocatalysts (Table S3); comparative table for photocatalytic H2O2production by various photocatalysts (Table S4) (PDF)

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#S.D. and S.P.T have equal contributions.

Notes

The authors declare no competing financial interest.

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Fault Diagnosis in Chemical Reactors with Data-Driven Methods

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ABSTRACT

This study investigates fault diagnosis, encompassing fault detection, isolation, and estimation, with experimental data in acontinuous stirred-tank reactor (CSTR) for the liquid-phase catalyticoxidation of 3-picoline with hydrogen peroxide. Two key faults were examined: coolant inlet temperature spikes (fault 1) and 3-picoline feedconcentration decreases (fault 2). Data-driven methods, including randomforest (RF) and k-nearest neighbors (KNN), successfully detected, isolated, and estimated faults under nominal conditions. However, bothdata-driven and model-based residual generators were disrupted by a shiftin the heat transfer coefficient (U). An isolation forest (IF) algorithm wasused for anomaly detection and model recalibration, restoring modelbased performance. Updated data sets enabled RF and KNN to adapteffectively, demonstrating their scalability and adaptability. Experimental results highlight the strengths of both methods, advocating for a combined framework for robust fault diagnosis.

1. INTRODUCTION

Fault diagnosis is a critical aspect of process safety in the chemical industry, where the reliability of operations candirectly influence productivity, economic performance, and,most importantly, safety.1,2 Chemical processes are highlycomplex and involve multiple interrelated variables, oftenoperating under extreme conditions of temperature, pressure, and chemical reactivity. Any undetected faults within thesesystems can lead to equipment malfunction, production losses, environmental hazards, and, in severe cases, catastrophicaccidents involving the loss of life. Therefore, ensuring robustand timely fault diagnosis systems is paramount to mitigaterisks and enhance the overall reliability of industrial operations.3

Fault diagnosis systems aim to detect, isolate, and estimate faults as early as possible to enable timely corrective actions, with each objective becoming progressively more challenging. In the chemical industry, faults can arise from various sourcesincluding sensor failures, actuator malfunctions, equipment degradation, and process disturbances. As processes becomemore automated and large-scale, human operators often facechallenges in manually identifying abnormal conditions due to the overwhelming amount of data generated by processsensors. Hence, automated fault diagnosis methods havegained prominence as key components of process safety frameworks. Early and accurate diagnosis not only prevents accidents but also optimizes maintenance schedules, reduces downtime, and enhances process efficiency. The necessity of robust fault diagnosis systems in the chemical industry has propelled the development of two primary approaches: model based and data-driven methods. 4–6

Each approach offers distinct advantages and faces its own set of challenges. Model-based fault diagnosis involvesconstructing a mathematical representation of the system tocompare real-time data with model-predicted outputs. Deviations from the expected behavior, known as residuals, are indicative of faults. 7 In contrast, data-driven approaches rely on historical operational data to train machine learning models to detect and classify faults. 8,9 While model-based methods leverage knowledge of the system's physics, datadriven techniques exploit patterns and trends hidden within large data sets. Model-based fault diagnosis techniques have long been favored in chemical process control due to their

strongfoundation in process physics. These methods are based onfirst principles such as mass and energy balances or empirical models that describe the behavior of the system. The core idea is to generate residuals, or error signals, bycomparing the system's actual measurements with predictions from the model. If a fault occurs, the residuals will deviate significantly from zero, indicating abnormal behavior. Popular model-based fault diagnosis techniques include observer based, 10–13 parameter estimation, 7,14 and parity space approaches. 15,16

Model-based approaches offer the advantage of physical insight, enabling fault detection, root cause analysis, and faultsize estimation without relying on extensive historical data. They are particularly reliable for rare or unmonitored faults butare heavily dependent on the accuracy of the system model. Developing such models for complex chemical processes is challenging and time-consuming, requiring a deep knowledgeof process dynamics and operational variables. Additionally, real-world systems evolve over time due to factors such as equipment aging or operational changes, leading to modelmismatch and potential diagnostic errors. Adaptive modeling isoften needed to recalibrate the model, adding complexity and effort, which limits the practical implementation of model based methods despite their reliability.

On the other hand, data-driven fault diagnosis methods have gained increasing popularity with the rise of machine learning and big data analytics. Data-driven approaches do not require explicit knowledge of the underlying system dynamics.8,17Instead, they rely on patterns, correlations, and anomaliespresent in the historical process data. Techniques such asrandom forest18,19 (RF), isolation forest20,21 (IF), supportvector machines 22,23 (SVM), artificial neural networks 24-26(ANN), and principal component analysis22,27,28 (PCA) are commonly used in data-driven fault diagnosis applications. John MacGregor and Nomikos have made monumental contributions to process monitoring and fault detection inchemical processes, proposing a PCA-based approach using only data from successful batches to monitor a styrene-butadiene semibatch reactor.29 An advanced fault isolationmethod has been developed to handle both simple and complex faults by extracting fault signatures and comparing them with a fault library of historical data. 30 Additionally, afault-tolerant control strategy employing datadriven latentvariable models constructed from historical process data ishighlighted, emphasizing their reduced dimensionality and interpretability. 31 Recent studies, particularly in Industrial & Engineering Chemistry Research (I&ECR), have focused onintegrating machine learning (ML) and artificial intelligence (AI) with process monitoring and fault diagnosis. 32 By linking fault diagnosis models with real-time digital replicas of physical systems, these approaches enable proactive maintenance, predictive analytics, and performance optimization.33–36The strength of data-driven methods lies in their ability tohandle complex, nonlinear systems without the need fordetailed models of the process. They are particularly wellsuited for systems in which developing an accurate model isimpractical or infeasible. Moreover, once trained, data-drivenmodels can be deployed to monitor systems in real-time and detect a wide range of faults with minimal human intervention. These methods are also highly scalable, making themapplicable to large and complex processes with multiplesensors and data points. However, the major challenge associated with data-driven approaches is the requirement for large and diverse data sets. In many cases, particularly for faultsize estimation, data-driven models often require training ondata sets that encompass a diverse range of fault scenarios ensure satisfactory performance. 37,38 In the chemical industry, where processes often run under nominal operating conditions for extended periods, it is difficult to obtain sufficient datarepresenting the various faulty states.

Recently, combined fault diagnosis methods, which incorporate the strengths of model-based and datadriven techniques, have emerged as a promising solution to overcome the limitations of both approaches.39–41 These methods aim to integrate the physical insights offered by modelbasedapproaches with the pattern-recognition capabilities of datadriven methods. For instance, model residuals can be used asinputs to a machine learning model, enabling more accurate fault classification. Alternatively, data-driven methods can be used to update model parameters in real-time, improving themodel's adaptability to changing system conditions.

Despite the potential advantages of combined approaches, their practical implementation remains limited, especially when applied to experimental data for achieving all three aspects of fault diagnosis: detection, isolation, and estimation. Moststudies in the literature focus on simulation-based validation, where fault scenarios can be artificially generated and tested. For example, the Tennessee Eastman (TE) process is acommon testbed for data-driven and model-based methods.42-44 As for experimental study, a hybrid approach combining an extended Kalman filter (EKF) with aprobabilistic neural network classifier has been successfully applied for fault detection and diagnosis in fed-batch and batchreactors, providing accurate monitoring through the estimation of reactor parameters and classification of fault types.45 Arobust fault detection methodology for hybrid process systems, incorporating tools from unknown input observer theory and Lyapunov stability, has been developed to reliably distinguishbetween faults, mode transitions, and uncertainties.46 A hybriddata/model-based approach combining SVM with an observeris proposed for fault detection and isolation in nonlinearchemical reactions, effectively reducing the reliance on precise process models or extensive training data.23 A hybrid modelcombining first-principles and neural networks was developed for automatic fault detection and identification, leveraging both simulation data and historical process information. Tested onreal data from a methanol-water distillation column, thismethod outperformed traditional first-principles models by effectively identifying faults and demonstrating its potential forapplication in refining and petrochemical processes.47

Recently, we introduced a comprehensive fault diagnosis methodology for a CSTR chemical reaction system, leveragingmodel-based residual generators as estimators, systematic dataprocessing to mitigate noise, and predefined thresholds forfault alarms. These residual generators, designed as functional observers decoupled from disturbances, estimate fault sizes. Fault isolation is achieved through multiple independent residual generators. 48,49 The experiment successfully demonstrated the effectiveness of fault diagnosis in a CSTR acrossvarious fault scenarios. 50 During these experiments, an intriguing phenomenon caught our attention: after switching equipment, the previous model's performance declined. Further experimentation suggested that this might be due to a change in the heat transfer coefficient. This raised important questions about how to detect model mismatches or parameter changes and how we could leverage known experimental datato enhance the fault diagnosis. Addressing these issues was essential to our work.

In this study, we examine data-driven approaches for fault diagnosis in chemical reactors, comparing them with modelbased observers. While model-based residual generators excelin robustness and accuracy, especially when the system'sdynamics are well understood, data-driven methods like RFand KNN offer promising, scalable solutions. Whereas both methods performed well under nominal conditions, systemparameter changes, like shifts in the heat transfer coefficient (U), posed challenges for both approaches. To address this, we implemented an isolation forest (IF) algorithm for anomalydetection and model recalibration. The study shows that combining data-driven and modelbased methods can enhancefault diagnosis, with data-driven techniques becoming more robust after training in updated system conditions.

2. BACKGROUND AND METHOD

This section introduces the reaction system and summarizes key results achieved using model-based fault diagnosistechniques. However, some challenges remain that cannot beeffectively addressed by model-based methods alone. This limitation motivates the exploration of a hybrid approach

that leverages the strengths of both model- and data-driventechniques. Additionally, the fundamentals of the data-drivenmethods applied in this work are briefly outlined.

2.1. Reactor (CSTR) Model and Experiment Setup.

The N-oxidation of alkylpyridine is a crucial reaction in drug synthesis and pharmaceutical applications. Studies have shownthat a continuous stirred tank reactor (CSTR) is an effective system for the N-oxidation of 3-picoline using hydrogen peroxide, 51 a process aligned with green chemistry principles. The reaction mechanism is illustrated in Figure 1.

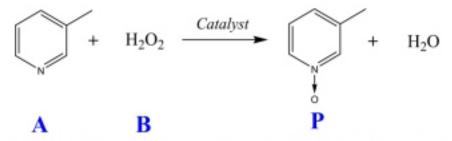


Figure 1. Reaction of catalyzed 3-picoline (A) N-oxidation with hydrogen peroxide (B) to produce N-oxidized 3-picoline (P) and water.

The CSTR setup and experimental process are schematically represented in Figure 2, with a 50 mL jacketed glass The objective is to detect, isolate, and estimate two faults in this process (for detailed information, refer to the cited work). During the experiments, faults were introduced, and sensorreadings (TT1 for reactor temperature, TT2 for jackettemperature, and AT for 3-picoline concentration) were used to successfully diagnose faults. This was achieved using model based residual generators derived from the functional observer applied to the system equations (eqs 1–5). Fault detection, isolation, and estimation were successfully achieved in our experiments using model-based residual generators.50 The reaction model is built with mass and energy balances, where CA is the concentration of reactant 3-picoline, CB is the concentration of reactant hydrogen peroxide, w(t) is theunknown kinetics variation, R(CA,CB,T) is the reaction rate, Tis the reactor temperature, $\delta t = 0.1$ s as the discretization time interval, and Tj is the jacket temperature. The system dynamics and the specifics of the residual generators are detailed as follows:50

$$\begin{split} C_A(k+1) &= C_A(k) + \delta_l \bigg(\frac{F}{V} (C_{A,\text{in}} - C_A(k) - f_2(k)) \\ &- (1 + w(k)) R(C_A(k), C_B(k), T(k)) \bigg) \end{aligned} \tag{1} \\ C_B(k+1) &= C_B(k) + \delta_l \bigg(\frac{F}{V} (C_{B,\text{in}} - C_B(k)) \\ &- (1 + w(k)) R(C_A(k), C_B(k), T(k)) \bigg) \end{aligned} \tag{2} \\ T(k+1) &= T(k) + \delta_l \bigg(\frac{F}{V} (T_{\text{in}} - T(k)) \\ &- \frac{UA}{\rho C_p V} (T(k) - T_j(k)) \bigg) \\ + \delta_l \bigg(\frac{-\Delta H}{\rho C_p} (1 + w(k)) R(C_A(k), C_B(k), T(k)) \bigg) \end{aligned} \tag{3} \\ T_j(k+1) &= T_j(k) + \delta_l \bigg(\frac{F_j}{V_j} (T_{j,\text{in}} - T_j(k) + f_1(k)) \\ &+ \frac{UA}{\rho_j C_{p,j} V_j} (T(k) - T_j(k)) \bigg) \end{aligned}$$

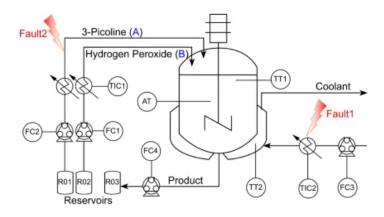


Figure 2. Schematic of experimental apparatus, indicating sensor and control locations.

$$y_1 = C_A(k) + \eta_1(k)$$

 $y_2 = T(k) + \eta_2(k)$
 $y_3 = T_j(k) + \eta_3(k)$ (5)

Residual generators were constructed for this system to implement fault diagnosis, which remain below a threshold ifthere is no fault happening and respond to faults (f1,f2), and toprovide an estimate of these fault sizes. For fault 1, a spike in coolant feed temperature Tj,in, thefollowing functional observer was built,

$$z_1(k + 1) = \alpha_1 z_1(k) - \frac{UA(1 - \alpha_1)}{\rho_j C_{p,j} F_j} y_2'(k)$$

 $- \left(\frac{(1 - \alpha_1)^2 V_j}{\delta_t F_j} - 1 + \alpha_1 - \frac{UA(1 - \alpha_1)}{\rho_j C_{p,j} F_j} \right) y_3'(k)$ (6)
 $r_1(k) = z_1(k) + \frac{(1 - \alpha_1) V_j}{\delta_t F_j} y_3'(k)$ (7)

Whereas the detection of fault 2, inlet feed ratio of $C_{A,in}$, and the 3-picoline concentration decrease, similarly, the following functional observer was built.

$$\begin{split} z_2(k+1) &= \alpha_2 z_2(k) + \left(\frac{(1-\alpha_2)^2 V}{\delta_t F} - 1 + \alpha_2 \right) y_1'(k) \\ &+ \frac{\rho C_p}{-\Delta H} \left(\frac{(1-\alpha_2)^2 V}{\delta_t F} - 1 + \alpha_2 \right) \\ &- \frac{UA(1-\alpha_2)}{\rho C_p F} \right) y_2'(k) + \frac{UA(1-\alpha_2)}{-\Delta HF} y_3'(k) \\ r_2(k) &= z_2(k) - \frac{(1-\alpha_1) V}{\delta_t F} y_1'(k) - \frac{(1-\alpha_1) V \rho C_p}{\delta_t F(-\Delta H)} y_2'(k) \end{split}$$
(8)

In eqs 6–9, y 1, y2, andy3represent the output measurements in deviation form, α 1 and α 2 are tunable parameters that represent the observer eigenvalues, r1 and r2 are the residuals that also represent estimates of f1 and f2, respectively, under the assumption that they are of step or ramp type. The details of individual parameter values are given in Table S1.50

It is important to emphasize that the residual generators are unaffected by fluctuations in the reaction kinetic rate, which could otherwise introduce significant errors. Moreover, the two residual generators are decoupled, allowing for effective faultisolation. The experimental results demonstrated successful fault diagnosis for both faults, providing highly accurate estimates of their magnitudes. 50

2.2. Random Forest (RF), K-Nearest Neighbors (KNN), and Artificial Neural Networks (ANN) for Fault Diagnosis. The random forest regressor is an ensemble learning technique tailored for regression tasks. 18 It constructs multiple decision trees during training and averages their predictions to enhance accuracy and mitigate overfitting. Unlike individual decision trees, which are prone to highvariance and overfitting, random forest uses bagging (bootstrapaggregation) and random feature selection to build a more robust and generalized model. This approach reduces the likelihood of overfitting and improves the model's performance on unseen data.

The random forest algorithm applies bootstrap sampling to generate multiple training subsets. Given a data set= $\{\}$ D=(x,y)IIi

n 1, where xi is the input feature and yi is the target variable, the algorithm generates t bootstrap samples. Abootstrap sample Db is used to train the decision tree. Theremaining data, called the out-of-bag (OOB) sample, can be used to estimate the model's performance. Each decision tree in the forest is constructed at each node; a subset of features $Fm \subseteq F$ is selected at random. The splitting criterion used in regression is typically based on minimizing the mean squared error (MSE). Once the tree is fully grown (or meets a stopping criterion, like maximum depth), it can make predictions for new data. The prediction for a new point x is the average of the target values y for all samples that fall into the same leaf node as y:

$$\hat{y}(x) = \frac{1}{N_{\text{leaf}}} \sum_{i \in \text{leaf}(x)} y_i$$
(10)

After multiple decision trees $\{T_1, T_2, ..., T_t\}$, the final prediction for new input x is obtained by averaging the predictions of all trees:

$$\hat{y}(x) = \frac{1}{t} \sum_{i=1}^{t} \hat{y}_{i}(x)$$
(11)

Where y(x) i is the prediction of the i-th tree.

The k-nearest neighbors (KNN) regressor is a nonparametric, instance-based learning algorithm used for regression tasks.52 Unlike statistical model-based methods, such as random forest, KNN does not require explicit trainingor model fitting. Instead, it predicts the target value byaveraging the target values of the k-nearest neighbors in thefeature space, relying solely on the stored training data. KNN uses a distance metric, typically the Euclidean distance, to identify the closest training points to a newquery point. In KNN regression, the predicted value for a given query point is the average of the target values of its k-nearestneighbors. The hyperparameter k determines how manyneighbors are considered in the prediction. Because KNN does not build a model, the computational cost of training isminimal, but predictions can be slower, especially with large data sets.

The algorithm for KNN can be explained as follows: for a given test point x, KNN calculates its distance to all trainingpoints using a predefined distance metric. The most commondistance metric used is the Euclidean distance. Once the distances between the test point x and all training points are computed, the algorithm identifies the k-nearest neighbors by selecting the k points with the smallest distances. Let

Nk(x) bethe set of k-nearest neighbors of x. Then, in KNN regression, the predicted value y(x) for a new point x is computed as theaverage of the target values of its k-nearest neighbors:

$$\hat{y}(x) = \frac{1}{k} \sum_{i \in N_k(x)} y_i$$
(12)

Artificial neural networks (ANNs) are a class of machine learning algorithms inspired by the structure and functioning biological neural systems.53 ANNs learn complex patterns from data through interconnected layers of nodes or "neurons." These algorithms are widely used in tasks such asclassification, regression, and generative modeling, when dealing with high-dimensional or unstructured data.

ANNs consist of an input layer, one or more hidden layers, and output layer. The hidden layers contain neurons that applylinear transformations followed by activation functions to theinput data, allowing the network to model nonlinearrelationships. The training process involves iteratively updatingthe network parameters (weights and biases) using optimization algorithms, e.g., stochastic gradient descent (SGD), tominimize a loss function, such as mean squared error (MSE) orcross-entropy loss. Nowadays, neural networks are the backbone of deep learning and a cornerstone of many stateof-the-art AI systems.54

The algorithm for training a neural network can be summarized as follows: given a data set of input—outputpairs (X,Y), the network predicts outputsyby applyingforward propagation through its layers. The error between the predicted and actual outputs is quantified using a loss function, and the gradients of this loss with respect to the networkparameters are computed using backpropagation. Finally, the are updated in the direction of the negative gradient by using an optimization algorithm. This process is repeated iteratively, until the model converges to an optimalset of parameters. The mathematical formulation for a single neuron in a neural network is where x is the input, W is the weight vector, b is the bias, z is the linear combination, and f(z) is the activation function. The activation function introduces nonlinearity, enabling the network to model complex patterns.

$$z = W^{\mathsf{T}}x + b, \ a = f(z) \tag{13}$$

2.3. Isolation Forest for Anomalies Detection.

Anomaly detection is a critical task across various domains, such as industrial monitoring, fraud detection, cybersecurity, and medical diagnostics. The objective is to identify datapoints that deviate significantly from expected patterns, whichmay indicate rare but important events such as systemmalfunctions, fraudulent activities, or network intrusions. This task is challenging due to the complexity and variability of real-world data, where normal behavior can fluctuate significantly and anomalies may be subtle or occur in highdimensional spaces. Isolation forest (IF), introduced by Liu et al.,55 offers a novelapproach to anomaly detection by focusing on the concept ofisolation rather than traditional distance or density-basedmeasures. The core principle of IF is that anomalies are "fewand different", making them easier to isolate from the rest of the data. Instead of evaluating a point's relative position within the data set, IF isolates each data point by recursively partitioning the data set through random splits. IF is an efficient algorithm for anomaly detection, focusing on isolating data points by recursively splitting the data set. Anomaliesstand out because they differ significantly from normal datapoints, and as a result, they are isolated more quickly during the partitioning process. The IF method offers several key advantages: (I) scalability: the algorithm scales linearly with the data set size, making itideal for large-scale applications; (ii) no assumptions aboutdata distribution: unlike methods that rely on specific datadistribution assumptions (e.g., Gaussian), IF is distribution agnostic, enhancing its robustness across various domains; (iii) handling high-dimensional data: IF performs effectively onhigh-dimensional data sets, avoiding the "curse of dimensionality" that hampers many traditional approaches.

The path length h(x) represents the number of edges traversed in an isolation tree before a point x is isolated. Anomalous points, which are more distinct, tend to have shorter path lengths. The anomaly score is based on this pathlength but normalized to fall within the range [-1, 1] where -1 indicates anomalies and 1 represents normal data points. Given a point x, its anomaly score s(x) is calculated as

$$s(x) = 2\left(\frac{E[h(x)]}{c(n)} - 0.5\right)$$
(14)

Where E[h(x)] is the average path length of point x across the isolation trees and c(n) is the normalization factor, representing the average path length for a normal point in adata set of size n. The IF algorithm operates as follows: data setX of size n and a number of trees t are subsampled with size ψ , forming a subsample $X\psi$. Then the subsample is recursively partitioned with a selected feature f and a split value p within the range of the feature, until each data point is isolated or the tree reaches a maximum depth L. For each point x, traverse the isolation tree to compute its path length h(x). Finally, calculate the anomaly score s(x), if it is positive, it indicates a normal point; on the contrary, if negative, it represents an anomaly.

3. RESULTS AND DISCUSSION

The sensor data sets were collected from open-loop experiments, as previously described and consistent with our earlier study.50 To reduce noise and ensure comparability, the datasets were normalized, and a 1200-point moving average filterwas applied to the data-driven fault prediction outputs, andresidual signals were filtered using fast Fourier transform(FFT). Data-driven training and analysis were conducted on aLenovo ThinkPad P53 (Intel 9850H CPU, Nvidia QuadroRTX 5000 mobile GPU, 128 GB RAM) using Python (Scikitlearn, TensorFlow, and PyTorch) and MATLAB on SlackwareLinux. Detailed parameters for RF, KNN, and IF are provided in Table S2 and ANN and RNN are provided in Table S3.

3.1. Data-Driven Fault Diagnosis. The residual generators have demonstrated their ability to detect, isolate, andestimate faults in the CSTR process. Following theexperimental runs, the collected data prompted the question of whether data-driven methods could also be applied for fault diagnosis. A total of 19 data sets with sensor readings and known fault sizes were used, with one serving as the test set and the remaining 18 as the training set. We applied methods: random forest (RF) and k-nearest neighbors (KNN). RF is a model-based statistical approach, while KNN isnonparametric, providing a representative comparison fortesting data-driven methods.

Since we already know the exact time and magnitude of the faults from the experimental data, training the data-drivenmodels does not require knowledge of the system equations-(eqs 1–5) or the model-based residual generators (eqs 6–9). Instead, the sensor readings are used to train the models solelyon the basis of the fault occurrence times and fault sizes, making this a "model-less" approach to fault diagnosis. The detection thresholds were determined using a Bayesian changepoint detection mechanism 56 with one data set and tested across all data sets, ensuring effective and accurate fault

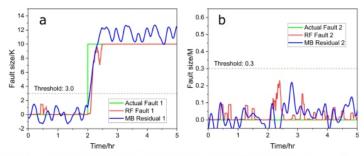


Figure 3. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with RF and compared to model-based (MB) residual generators (blue), under scenario a.

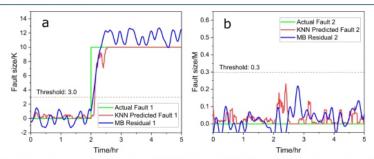


Figure 4. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with KNN and compared to model-based (MB) residual generators (blue), under scenario a.

detection. Additionally, the data-driven regressors accurately converged to the actual fault size, demonstrating a strong performance in fault estimation.

A comprehensive analysis was conducted as follows:

a. Only one fault is happening,

$$f_{\rm l}(t) = \begin{cases} 0, & t < 2 \, {\rm h} \\ 10, & t \geq 2 \, {\rm h} \end{cases}, \quad f_{\rm l}(t) = \begin{cases} 0, & t < 2 \, {\rm h} \\ 0, & t \geq 2 \, {\rm h} \end{cases}$$

With only one fault introduced into the system fault 1, caused by a spike in coolant inlet temperature the expected outcome is accurate detection and estimation of the fault size, while predictions for fault 2 remain low, as no fault associated with 3-picoline concentration is present.

As shown in Figure 3a, fault 1 is introduced at the 2 h mark, at which point the RF regressor (red line, eq 11) promptly detects the fault (represented by the green line) and responds within 5 min, converging to the correct fault size of 10 within 30 min. A threshold of 3 could be applied for fault detection, when the signals of the residual generator or RF regressorexceed this threshold, a fault is then alarmed. We also evaluated the performance of the moving average filter, with a 1200-point window proving to be an effective choice, with R2 = 0.96, and the results are provided in Table S3. Theperformance of RF is comparable to that of a model-basedresidual generator (blue line, "MB" for model-based) with aslight delay, demonstrating the efficiency of RF in faultdetection and size estimation. In Figure 3b, the signal for fault2 remains low throughout the test, confirming that no fault related to the decrease in 3-picoline concentration is present, of which the residual generator (blue line) provides a similar result. Similarly, a threshold could be set at 0.3, and both signals remain below this threshold, indicating that there is no fault happening. This accurate isolation further validates therobustness of the RF model for single-fault scenarios. These results highlight the ability of data-driven methods to both detect and isolate faults with precision comparable to modelbased residual generators, even in complex system dynamics. The same procedure was applied using the k-nearest-neighborhood (KNN) regressor, and the results are shown in Figure 4. In Figure 4a, the KNN model (red line, eq 12) swiftly responds to the introduced fault 1, similar to the RF results, detecting the fault promptly and converging toward the correct fault size. In Figure 4b, the signal for fault 2 remains consistently low throughout, indicating successful fault isolation and confirming that no fault is associated with the decrease in the 3-picoline concentration. These results demonstrate that the KNN regressor, like the RF model, achieves both fault isolation and size estimation effectively, reinforcing the viability of data-driven approaches in fault diagnosis tasks. In both methods, there is a slight delay of datadriven methods compared with model-based residual generators, showing a possible quicker response in residual signals.b. Only one fault is happening,

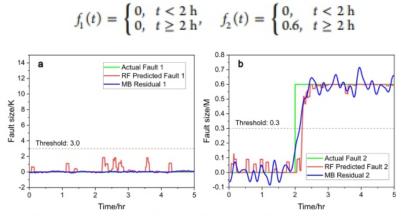


Figure 5. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with RF and compared to model-based (MB) residual generators (blue), under scenario b.

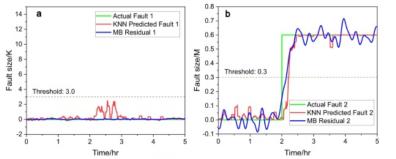


Figure 6. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with KNN and compared to model-based (MB) residual generators (blue), under scenario b.

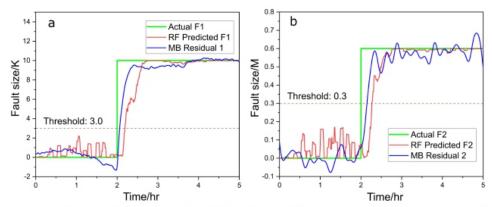


Figure 7. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with RF and compared to model-based (MB) residual generators (blue), under scenario c.

In the reversed scenario, where the only fault occurring in the system is fault 2�a decrease in the 3-picoline feed inletconcentration the expectation is that the model will detectand estimate the size of fault 2, while the prediction for fault 1 remains low, as no coolant inlet temperature spike is present. In Figure 5a, the fault 1 prediction remains consistently lower than 3, indicating no fault related to the coolant inlettemperature, as expected. In Figure 5b, the model responds to fault 2 swiftly and converges

to the correct fault size within 30min, showing promising results in fault detection and estimation. Similarly, when the KNN regressor is tested in

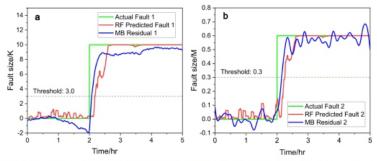


Figure 8. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with KNN and compared to model-based (MB) residual generators (blue), under scenario c.

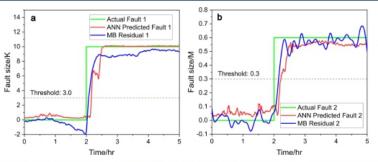


Figure 9. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with ANN and compared to model-based (MB) residual generators (blue), under scenario c.

this scenario: in Figure 6a, the predicted value for fault 1 stays low, though slightly higher than that of the RF model, stillconfirming no fault. In Figure 6b, the KNN model promptlydetects fault 2, accurately estimating its size with a short response time.

Despite minor discrepancies observed in both fault scenarios, the models consistently detect faults when theyoccur, although the precision of the magnitude estimation varies. Crucially, the timing of fault detection is accurate inboth cases, with the models successfully identifying both theonset and the resolution of faults. This is especially clear in the case of fault 2, where the predicted and actual fault sizes nearly converge during the period of sustained fault, demonstrating the models' effectiveness in accurately tracking fault behavior over time. For both RF and KNN, a slight delay is also present compared with residual signals.

c. Two faults are happening,

$$f_1(t) = \begin{cases} 0, & t < 2 \text{ h} \\ 10, & t \geq 2 \text{ h} \end{cases}, \quad f_2(t) = \begin{cases} 0, & t < 2 \text{ h} \\ 0.6, & t \geq 2 \text{ h} \end{cases}$$

Figures 7 and 8 present the comparison between the actual and predicted fault sizes for fault 1 and fault 2, using twodifferent machine learning models: RF and KNN, respectively. Each figure consists of two subplots: Figures 7a and 8 are present fault 1, and Figures 7b and 8b represent fault 2. Inboth cases, the predictions are compared against the ground offault sizes, along with model-based residual signals.

The random forest model shows a very close alignment between the predicted fault size (red line) and the actual faultsize (green line). The model captures the onset of the faultaround 2 h, and the prediction remains accurate throughoutthe fault duration. As for the residual signal, there is a smalldeviation, less than 1.0 K, especially during the sustained faultevent (fault size ~10), which is acceptable. The sharp rise andstable fault size during this period suggest that random foresthandles significant faults efficiently and with high precision. Prior to the major fault event, the RF model exhibits somesmall fluctuations in the predicted fault size that occur in the time range between 0 and 1.5 h. These small

deviations do not affect the overall value as they are below the detection threshold.

The KNN also shows strong predictive performance with the predicted fault size (red line) following the actual fault size(green line) very closely during the main fault period. The riseof the fault at approximately 2 h and the sustained fault are captured well.

For fault 2, both models exhibit similar performance, but the random forest model shows more variability and fluctuation inits predictions. This might indicate that random forest is moreprone to overfitting to noise in cases where the fault dynamics more complex, but all fluctuations are well below detectionthresholds. As for the model-based residual signal (MB, theblue line), the noise level is much higher than data-driven

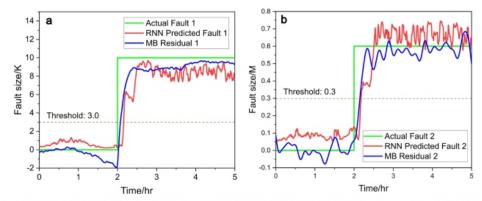


Figure 10. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with RNN and compared to model-based (MB) residual generators (blue), under scenario c.

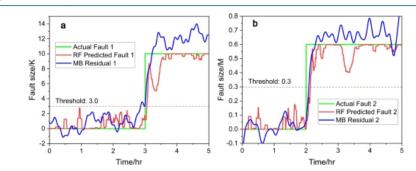


Figure 11. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with RF and compared to model-based (MB) residual generators (blue), under scenario d.

methods; however, it has a better prompt response to the fault happening.

Neural networks have gained popularity in recent years to their ability to model complex, nonlinear systems effectively. Building on the fault diagnosis conducted using KNN and RF, we also applied artificial neural networks (ANN) and recurrentneural networks (RNN) to explore their effectiveness. Thesemethods were chosen to leverage their capacity for capturing intricate patterns and, in the case of RNN, for addressing temporal dependencies within the data. The fault diagnosis achieved using an ANN and an RNN is illustrated in Figures 9 and 10. In Figures 9a and 10a, the ANN and RNN successfully identify fault 1 by comparing the actual fault (green line) with the predicted fault (red line). The threshold of 3.0 (dotted line) is used to detect fault occurrences. Similarly, in Figures 9b and 10b, fault 2 is diagnosed with a threshold of 0.3, where the ANN and RNN both capture the fault dynamics through the alignment of the actual fault (green line) and predicted fault (red demonstrates the effectiveness of neural networks in detecting and predicting fault conditions. Figure 10 plotted using RNN appears to exhibit higher noise levels in the predicted compared with Figure 9. This increased noise could be attributed to RNN's sensitivity to temporal dependencies, which may amplify variations in the data. While the neural network demonstrates strong performance in fault diagnosis, the computational time

required is much higher compared to RF and KNN. The average timeconsumption for a single cycle of training and testing using KNN, RF, ANN, and RNN is detailed in Table S5. Additionally, Table S6 compares the performance of these methods in estimating fault sizes against the actual fault values. We will present results only for KNN and RNN, as they demonstrate satisfactory performance. d. Two faults happening at different times

$$f_1(t) = \begin{cases} 0, & t < 3 \text{ h} \\ 10, & t \ge 3 \text{ h} \end{cases}, \quad f_2(t) = \begin{cases} 0, & t < 2 \text{ h} \\ 0.6, & t \ge 2 \text{ h} \end{cases}$$

It is crucial to verify fault isolation with experimental data that fault 1 and fault 2 are happening at two different times. The RF model accurately predicts the onset and magnitude of fault 1 in Figure 11a. The predicted fault size (red line) closely follows the actual fault size (green line), particularly during the critical period after 3 h, where the fault size rises sharply and stabilizes around a fault size of 10. The model consistently captures the duration and magnitude, demonstrating a high prediction precision during the main fault event.

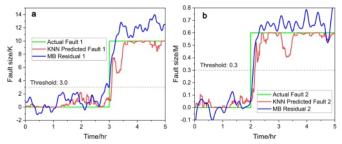


Figure 12. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with KNN and compared to model-based (MB) residual generators (blue), under scenario d.

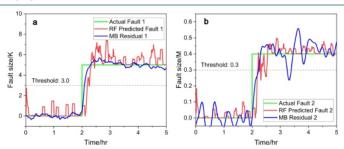


Figure 13. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with RF and compared to model-based (MB) residual generators (blue), under scenario e.

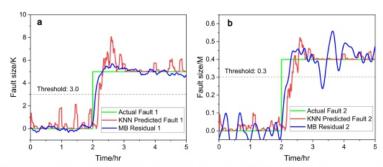


Figure 14. Predicted fault size (red) vs actual fault size (green), (a) for fault 1 and (b) for fault 2, with KNN and compared to model-based (MB) residual generators (blue) under scenario e.

The KNN model similarly demonstrates strong predictive performance for fault 1 in Figure 12a, with the predicted faultsize (red line) following the actual fault size (green line)closely. The model effectively captures the sharp rise in the fault size around 3 h, maintaining accuracy throughout the sustained fault period. As for the model-based residual signal (MB, the blue line), we observe an overshoot, giving a 1.2

Koverestimation of fault size, which is within the experimental tolerance.

For fault 2, in Figure 11b, the predicted random forest model fault size (red line) captures the general trend of thefault, which begins around 2 h and persists until the end of thetest. After the fault reaches its maximum (fault size 0.6), therandom forest model stabilizes and follows the actual fault size (green line) more closely. In Figure 12b, the KNN modeloffers slightly smoother predictions for fault 2. While the predicted fault size (red line) does fluctuate in the early fault period (around 3 h), the overall estimation is sound.

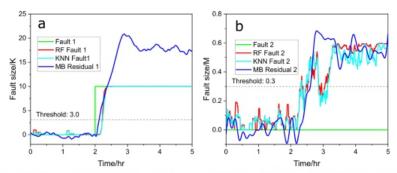


Figure 15. Predicted fault size vs actual fault size with residual generators, RF (red) and KNN (cyan), (a) for fault 1 and (b) for fault 2, and compared to model-based residual generators (blue).

These results exhibit successful fault isolation; the noise level is not too high to invalidate fault detection, and especially theestimation for both fault sizes is remarkably good. The slowerresponse compared with residual generators is not unforeseeable, as model-based methods have the edge in the system dynamics.

e. Two smaller faults are happening

$$f_1(t) = \begin{cases} 0, & t < 2 \text{ h} \\ 5, & t \ge 2 \text{ h} \end{cases}, \quad f_2(t) = \begin{cases} 0, & t < 2 \text{ h} \\ 0.4, & t \ge 2 \text{ h} \end{cases}$$

Figures 13 and 14 demonstrate the performance of fault prediction models RF and KNN for a smaller fault scenario, respectively. Each figure consists of two subplots: a representsfault 1 and b represents fault 2, where the faults are smaller inmagnitude compared to previous scenarios.

In this smaller fault scenario, the random forest model shows relatively good performance in Figure 13a. The predicted faultsize (red line) aligns with the actual fault size (green line), particularly after the fault onset around 2 h. However, compared to the larger fault scenario, the random forestmodel introduces more variability in its predictions. The predicted fault size fluctuates around the actual fault size, especially after 2.5 h, where the model tends to overestimate the fault magnitude. The KNN model provides more stablepredictions for fault 1 under the smaller fault scenario in Figure 14a after 3 h. The predicted fault size (red line) follows the actual fault size (green line) closely, with fewer deviations compared with the random forest model. There is a minoroverestimation of the fault size around 2.5 h, but overall, the model tracks the actual fault size more consistently. The residual signal shows superb performance with both accuracy and responsiveness.

For fault 2, the random forest model (Figure 13b) and the KNN model (Figure 14b) exhibit a slight overestimate at the beginning of the fault occurrence, and the KNN model shows a larger overshoot at around 2.5 h. Overall, the KNN model's ability to provide smoother predictions with fewer oscillations makes it better suited for capturing the smaller faults in fault 2, as it appears less sensitive to noise or minor deviations in the data especially after 3 h.

Estimating smaller faults is generally more challenging due to the lower signal-to-noise ratio, a difficulty also withmodel-based observers, as discussed in our previous study.50The results demonstrate that data-driven methods (RF andKNN) are capable of diagnosing these minor faults effectively.

In conclusion, both models are effective at detecting smaller faults. Model performance: both random forest (RF) and knearest neighbors (KNN) demonstrated strong performance indetecting and isolating faults as well as estimating their sizes. The circular iteration ensured robustness and consistency across the data set. However, in all the cases, model-basedresidual signals show better responsiveness and faster fault detection. Accuracy: the average accuracy across all tests washigh, with minimal variance between data sets. This indicates the effectiveness of both RF and KNN across different operational conditions. Fault Isolation: both methods successfully isolated the faults in the system, identifying the affected components without false positives or significant misclassifications. Fault size estimation: the size estimation of faults was within acceptable error margins, showing that data-driven methods, when properly trained, can provide reliable fault magnitude estimates without relying on model-based equations.

3.2. Isolation Forest for Anomalies Detection. A major limitation of the model-based approach is that not all processescan be easily modeled, and even with a validated model, operational parameter changes may still occur. Similarly, datadriven methods cannot reliably diagnose faults without priorknowledge of faulty data sets or at least similar conditions fortraining. For example, during experimentation, it was observed that after reassembling the equipment the previously effective model-based residual generators showed reduced performance, while data-driven methods produced false detections. This setup considered a scenario in which only a single fault was present in the process, as described below.

$$f_{_{\! 1}}(t) = \left\{ \begin{matrix} 0, & t < 2 \text{ h} \\ 10, & t \geq 2 \text{ h} \end{matrix}, \quad f_{_{\! 2}}(t) = \left\{ \begin{matrix} 0, & t < 2 \text{ h} \\ 0, & t \geq 2 \text{ h} \end{matrix} \right.$$

The model-based residual signals are shown as blue lines. For fault 1 in Figure 15a, the model-based residual signalspikes above 20 before stabilizing around 17, which is over60% higher than the actual fault size, indicating a clearoverestimation. In Figure 15b, despite no fault being presentfor fault 2, the residual signal incorrectly estimates around 0.5, as a false positive. For the data-driven methods, RF (red lines) and KNN (cyan lines), the performance for fault 1 in Figure 15a is notably strong, accurately detecting and estimating thefault size at 10. However, both methods also produced falsepositives for fault 2 in Figure 15b. All methods failed under the system parameter change, which is expected. Data-driven methods struggled due to the lack of similar scenarios in thetraining sets, and the model-based observer failed because theparameter change caused a mismatch in the system model, invalidating the residual signals.

This was traced back to a change in the heat transfer coefficient with additional experiments. It is crucial to detectchanges in the heat transfer coefficient (U) before any faultsare introduced into the system. Early detection allows fortimely adjustments to the model parameters, preventingpotential losses due to detection delays when actual faultsoccur. Therefore, the goal is to detect shifts in U during thenonfault steady-state operation, enabling proactive model adjustments before faults impact system performance. In data-driven methods, detecting changes in the heat transfer coefficient (U) is equivalent to identifying anomalousdata sets compared to a set of nominal runs. Therefore, it isnatural to apply an anomaly detection algorithm, with Isolationforest (IF) being an ideal choice due to its simplicity and lowcomputational resource requirements. The isolation forest (IF)parameter settings include 200 estimators, a contaminationlevel of 0.2, and a maximum sample size of 256. Further detailson all parameter settings are provided in Table S2. The sensordata sets, collected before any faults were introduced into the system and with a known $U = 18 \text{ W/(m} 2 \cdot \text{K})$ (calculated using a heat transfer area of 0.08 m2), are used to train the IF model. After training, several data sets with potentially altered Uvalues are tested against the model, and anomaly scores are calculated according to eq 14. The results are summarized in Figure 16, where four datasets were randomly selected for testing: two with old U = 18

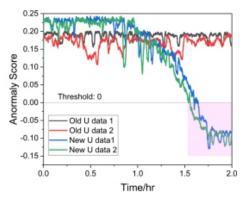


Figure 16. Isolation forest anomaly score for process U change.

 $W/(m2 \cdot K)$ and two with new U. In the figure, the two data sets with U=18 show consistently positive anomaly scores, indicating that the sensor readings align with the training set. In contrast, the two data sets with altered U values show agradual detection of anomalies by the IF model, with anomaly scores dropping sharply into the negative range and stabilizing around -0.1, as marked by the pink area in Figure 16. This successful anomaly detection confirms a change in U within these two data sets.

Once an anomaly is detected, the next step is to adjust the model parameter (U) to ensure that the residual generators continue to function optimally for fault diagnosis. There are numerous regressors available in the data-driven toolbox forestimation tasks. However, a key drawback of data-driven methods is that models must be trained with data sets corresponding to various known U values. In practice, the Uvalues may not always be known a priori, which can result inunreliable output estimates. To address this, we refer back to the system model described in eqs 1–5.

Assuming the system is in a steady state and no faults are present, we can derive an equation from eqs 1 and 3, resultingin eq 15. Additionally, by considering eq 4 alone, we can stablish eq 16. Both of these equations provide estimates of the U value, offering a more reliable approach than purely datadriven methods.

$$U_{1} = \frac{F}{VA(T(k) - T_{j}(k))} (T_{in} - T(k)) + \frac{-\Delta HF}{(T(k) - T_{j}(k))}$$

$$(C_{A,in} - C_{A}(k)) \qquad (15)$$

$$U_{2} = -\frac{F_{j}\rho_{j}C_{pj}}{(T(k) - T_{j}(k))} (T_{j,in} - T_{j}(k)) \qquad (16)$$

Direct calculation using eq 16, as shown in Figure 17a, incorporates a moving average filter with a window size of 6000 points. However, the results are suboptimal due to the high noise levels. Not all data sets can be reliably calculated using eqs 15 and 16 because the divisor involves temperature measurements, which are subject to sensor noise. As noted in eq 5, floating-point division can lead to significant errors when the divisor is small.

The mean values of the calculated lines in Figure 17a are 17.6 and 19.07, which are close to the actual value of $18 \text{ W/(m2 \cdot K)}$. However, the high noise level makes direct calculationimpractical for all data sets. This issue arises because thereactor jacket's thickness is only around 0.2 mm, resulting in alarge heat transfer coefficient. Consequently, the temperaturedifference between T-Tj is small and is exacerbated evenmore, as the measurement noises (eq 5) of temperature have avariance over 0.6, which can lead to significant errors in the calculation due to the small magnitude of the temperature difference. A more effective approach is to frame the problem as an optimization task. By rearranging eqs 15 and 16, we derive the residuals as follows:

$$Residual_1 = F\rho C_p(T_{in} - T) + -\Delta HF(C_{A,in} - C_A)$$

$$- U_1A(T - T_j) \qquad (17)$$

$$Residual_2 = -F_{\rho_j}C_{p,j}(T_{in} - T) - U_2A(T - T_j)$$
(18)

The objective function to be minimized is the sum of the squared residuals:

Objective =
$$\sum \text{Residual}_1(i)^2 + \sum \text{Residual}_2(i)^2$$
 (19)

under the condition with the lower bound $[U1\ U2] = [0\ 0]$. The optimization problem is solved using the global trustregion reflective algorithm, and the results are plotted in Figure 17b. The calculated curves exhibit much smoother behavior. The nominal runs, old U data 1 (black line) and old U data 2 (red line), fluctuate between 15 and 20, indicating stable butslightly lower U estimates over the actual value of 18. Theoverall trend for these nominal conditions remains relatively constant.

In contrast, the changed conditions, new U data 1 (green line) and new U data 2 (blue line), show fluctuations around

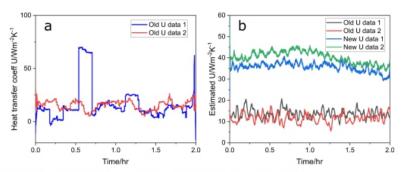


Figure 17. Direct calculation of U values for two nominal runs (a) and optimization estimates (b).

40 W/(m2·K), representing significantly higher U values compared to the nominal runs. This stark difference in Uvalues between the old U and new U conditions reflects achange in the heat transfer coefficient. The mean values of Figure 17b, calculated U, are summarized in Table 1.

Table 1. Calculated U Values from Optimization

Parameter	Old U	Old U	New U	New U
	data 1	data 2	data 1	data 2
$U/W/(m^2K)$	16.57	17.74	35.6	33.1

Using this method, following data triage, the experimental data revealed two distinct categories of U values: one centeredaround 18 W/(m2·K) (average 17.8, standard deviation 1.12) and another centered around 40 W/(m2·K) (average 38.9,

standard deviation 3.1), demonstrating good stability. After isolation forest (IF) anomaly detection, if a data set is flaggedas negative, the U value is recalculated to update the modelbased residual generators for fault diagnosis. The results are depicted in Figure 18, showing the faultdiagnosis for both fault 1 and fault 2 using parameter-updatedmodel-based residual generators. In Figure 18a on the left, theblue line represents the updated residual signal, which slightlyoverestimates the fault size after the 2 h mark, stabilizing around 11.5, within acceptable experimental tolerance. The redline shows the previous residual signal, as seen in Figure 15a, which deviates significantly from actual fault 1 (green line). In Figure 18b on the right, the results for fault 2 detection are shown, with no actual fault present (as indicated by the green line staying flat at zero). The updated residual signal (blueline), although showing some fluctuations, does not exceed the fault detection threshold. This updated signal performs

muchbetter than the previous residual signal in Figure 15b, indicatedby the red line. The procedure is outlined in the flowchart presented in Figure 19. Initially, sensor readings are analyzed using theisolation forest (IF) algorithm for anomaly detection. If no anomalies are found, the system proceeds with the modelbased residual generators using the default U value. However, if the anomaly scores drop below zero, indicating a potentialissue, then an optimization algorithm is triggered to recalibrate the U value. This updated U value is then used to adjust theresidual generators, allowing the model-based fault diagnosis tocontinue accurately under the new conditions. To apply a data-driven method, a backup model-based residual generator must be used for fault diagnosis in the newparameter-changed scenario until sufficient data sets with thenew parameter and fault conditions are collected. Once thesedata sets are validated against actual faults, they are incorporated into the data-driven RF and KNN training sets, labeled with the corresponding U value. This process updates the training sets, enabling data-driven methods to adapt to thenew conditions. To evaluate the efficacy of this approach, seven data sets with U = $40 \text{ W/(m2 \cdot K)}$ were used. Figure 20 presents the fault diagnosis results with RF and KNN. The same data set from Figure 13 is used as the testing set, and six new data sets

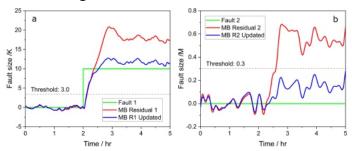


Figure 18. Model-based residual signal responses after parameter update (blue) and before update (red) compared with actual fault size (green).

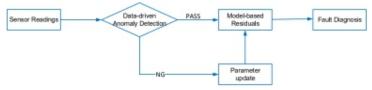


Figure 19. Flowchart of parameter update for model-based fault diagnosis.

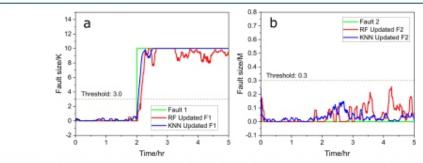


Figure 20. Predicted fault size vs actual fault size (green) with RF (red) and KNN (blue), (a) for fault 1 and (b) for fault 2 after training sets updated.

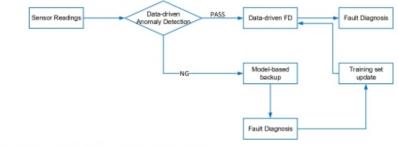


Figure 21. Flowchart of training set update for data-driven fault diagnosis.

marked with U = 40 were incorporated into the data-driven training set. This process allowed for data-driven approaches totrain and gain new knowledge about the system, thusenhancing fault diagnosis.

The result is shown in Figure 20, presenting the fault diagnosis results for both fault 1 and fault 2 using trainingupdated RF and KNN. In Figure 20a on the left, the red lineshows the RF model's prediction, which closely follows theactual fault size but fluctuates slightly after the initial detection, oscillating just below the true fault magnitude. The blue line, representing the KNN prediction, also aligns well with theactual fault, responding quickly and providing accurate stimates with fewer fluctuations compared with the RF prediction. In Figure 20b on the right, the results for fault 2 detection are displayed, where no actual fault is present (asindicated by the green line remaining flat at zero). Both the RF(red line) and KNN (blue line) predictions show minorfluctuations hovering slightly above zero. Nevertheless, these signals remain below the threshold of 0.3, indicating that there is no fault. As more training scenarios are gathered, it has been demonstrated that data-driven methods can also effectively perform fault diagnosis, matching the reliability of model-based approaches. Overall, the data-driven approaches (RF and KNN) demonstrate strong performance for fault 1, accurately detecting and estimating the fault size. In the case of fault 2, despite some fluctuations, there are no false positives andreliable fault isolation. This procedure is illustrated in Figure 21. Sensor readings are first fed into the anomaly detection process using theisolation forest (IF). If no anomaly is detected, the systemapplies the data-driven fault diagnosis as outlined in Section3.1. However, if an anomaly is detected, a backup model-basedfault diagnosis, as shown in Figure 19, is initiated. Oncesufficient new data with different fault scenarios are collected, the data-driven fault diagnosis model is retrained with the

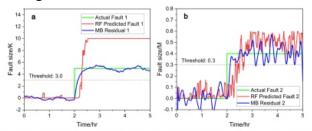


Figure 22. Predicted fault size vs actual fault size (green), (a) for fault 1 and (b) for fault 2 with RF (Red) and model-based residual generator (blue)

updated training set to enhance its fault diagnosis capabilities under the new conditions.

3.3. Challenges in Data-Driven Fault Diagnosis. Datadriven methods have shown great promise in fault diagnosis due to their ability to model complex systems using historicaldata. However, their effectiveness is highly dependent on thequality and diversity of the training data. Without a diversedata set that includes faulty scenarios across a wide range of operational conditions, these models often fail to generalize, leading to poor performance in real-world applications, asshown in Figure 15. When properly trained with comprehensive data that captures all potential fault types and operating conditions, data-driven models excel in faultdetection and diagnosis, offering accurate solutions, as shownin Figure 20. The key lies in ensuring that the training datacover all critical scenarios. The question arises whether data are required for all faultsizes. To explore this, a test was conducted using a training setwith only larger fault sizes (fault 1 at 10 and fault 2 at 0.6), while the testing set comprised smaller faults (fault 1 at 5 and fault 2 at 0.4). Figure 22 illustrates the fault diagnosis results for both fault 1 and fault 2, comparing the actual fault values, predicted fault values, and residual signals. This test provides insight into how well the models generalize to smaller faultsizes when they are trained only on larger faults. In Figure 22b, fault 2 is introduced at the 2 h mark, with thegreen line indicating the actual fault size of 0.6. Prior to this, the noise in both the predicted fault signal (red line) and theresidual signal (blue line) stays well below the threshold of 0.3, indicating no false positives before the fault occurs. After 2 h, both the predicted fault signal and the residual signal increase, with the red line overestimating the actual fault size of 0.4, while the blue residual signal fluctuates but stays close to the expected value of 0.4. Similar results were achieved using KNN; additionally, ANN and RNN were applied, yielding comparable results (not included in the article).

This phenomenon highlights the importance of having diverse training scenarios for accurate fault diagnosis, arequirement that may be difficult to achieve in real industrial settings. The question of what constitutes the minimal trainingset catches our attention. Further studies are needed to address this.

4. CONCLUSION

In this study, the effectiveness of data-driven methods such as RF and KNN for fault diagnosis in a CSTR chemical reactorsystem is evaluated compared to model-based residualgenerators. The focus is on detecting, isolating, and estimatingthe size of two key faults: fault 1, a coolant inlet temperaturespike, and fault 2, a decrease in the 3-picoline feedconcentration. Both RF and KNN demonstrated strongperformance under nominal conditions, accurately identifying faults and estimating their sizes. However, both data-drivenand model-based approaches faced difficulties after a change in the heat transfer coefficient (U), a process condition change, which led to misaligned predictions and false positives. Toaddress this, an IF algorithm is employed for anomaly detection, allowing system model recalibration and restoring the accuracy of the model-based residual generators. For datadriven methods, the inclusion of new data sets with updated parameters in the training successfully restored their performance.

Overall, while model-based methods remain reliable due to their deep understanding of system dynamics, data-drivenapproaches offer scalability and efficacy without the need fordetailed system models. The integration of both methods into a combined framework offers an optimal solution.

A challenge with data-driven methods is their reliance on diverse training data sets; without sufficient variety, faultdiagnosis accuracy diminishes, as seen in overestimations whentrained only on larger fault sizes. Future work will focus onenhancing the robustness of data-driven methods, particularlyby optimizing training sets with comprehensive fault scenariosand improving their ability to handle system parameterchanges to achieve fault diagnosis accuracy comparable tomodel-based approaches. And we will develop and test newmodels to enhance sensitivity and accuracy in detecting smallerfaults. The sensitivity analysis of heat transfer efficiency will beconducted when sufficient data covering a broader range of Uvalues becomes available. To further enhance the applicability of our approach, future work will explore strategies forreducing prediction delays, such as optimizing modelarchitectures for efficiency or implementing online learningtechniques to enable real-time performance. We will alsointegrate data-driven and model-based methods throughadaptive model adjustment, where data-driven techniques dynamically refine model parameters, and physics-guidedlearning, which implants physical constraints into neural networks. Additionally, transfer learning using model-based simulations will be explored to improve data-driven models inscenarios with limited real-world fault data.

■ ASSOCIATED CONTENT

*si Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.iecr.4c04042.Design parameters of CSTR experiment in Table S1.Parameters for all data-driven methods in Table S2 and Table S4. Moving average filtering performance in terms of determination coefficient in Table S3. Average time consumption for 1 cycle of training and testing in Table S5 and the performance across scenario a to e in Table S6 (PDF)

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Notes

The authors declare no competing financial interest.

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Celebrating the Birth Centenary of Quantum Mechanics: A Historical Perspective

Venkat Venkatasubramanian*

ABSTRACT

In July 1925, Werner Heisenberg submitted a paper to Zeitschrift f ür Physik entitled 'On quantum theoreticalreinterpretation of kinematic and mechanical relationships', thusgiving birth to quantum mechanics. In the following year, buildingon de Broglie's wave-particle duality, Erwin Schrödinger developedwave mechanics, and soon, Max Born provided a probabilisticinterpretation of the wave function. The theory was furtherenriched by the exclusion principle of Wolfgang Pauli and theuncertainty principle of Heisenberg, which ultimately led to thedevelopment of relativistic quantum mechanics by Dirac. TheCopenhagen Interpretation created a probabilistic framework forunderstanding the theory. Over the past century, quantummechanics has paved the way for advances in quantum fieldtheory, computing, and modern technologies. This historical narrative provides insights into the complex discovery process that led to the development of quantum mechanics, which can potentially guide novel breakthroughs amid challenging conceptual struggles, as seen in the field of artificial intelligence today.

1. WHY IS THE HISTORY OF QUANTUM MECHANICS IMPORTANT?

This year, we mark a historic occasion: the centenary of the birth of quantum mechanics. A hundred years ago, the field of quantum mechanics emerged through the pioneering efforts of primarily Werner Heisenberg, Erwin Schrödinger, and MaxBorn.1 Their groundbreaking contributions unraveled themysteries of the atomic world, transforming our understanding of reality itself. Quantum mechanics stands as one of humanity's most profound intellectual achievements. Its founding principles quantization, wave-particle duality, probability, uncertainty, and superposition dramatically redefined our understanding of the universe. Jagdish Mehra, the authoritative chronicler of the history of quantum theory, declared:1,2"The birth of quantum mechanics presents us with one of the most remarkable episodes in the history of science; it is as rich, complex, dramatic, and touching as any in the history of human thought."Quantum mechanics, with its further development asquantum field theory, is a magnificently beautiful theory, perhaps second only in its beauty to the general theory of relativity. But quantum mechanics is far more surprising thangeneral relativity in its strangeness with concepts such asquantum entanglement — "spooky action-at-a-distance," as Einstein put it-, that whisper yet-to-be-revealed deeper secrets of reality that seem almost mystical.

In the July 1925 paper,3 Heisenberg introduced matrix mechanics, marking the first comprehensive formulation of quantum theory that focused on observable quantities such as energy and spectral transitions. Between November 1925 and January 1926, Erwin Schrödinger developed wave presenting the now-famous Schrödinger equation, an alternative, yet equivalent, description of quantum systems. Soon, Max Born provided the probabilistic interpretation of thewave function, reshaping our notion of determinism and causality in physics.

These milestones were part of an extraordinary period of intellectual explosion, during which luminaries such as NielsBohr, Paul Dirac, and Wolfgang Pauli contributed to theframework that continues to underpin modern physics.1,2 This centennial is an opportunity to reflect on the profound human capacity for imagination and discovery. It is also an opportunity to marvel at "the unreasonable effectiveness of mathematics in the natural sciences," as Eugene Wigner wondered!4

This year, let us take a moment to appreciate and enjoy this crowning achievement of the human mind, which reveals themagnificent beauty of the hidden order of the cosmos. Thispaper is written with this objective in mind. Another importantreason for studying the early history of quantum mechanics is that it is one of the rare occasions when a considerable wealth offirst-hand accounts of momentous discoveries is available. Fortunately, many of the original architects of quantum theorylived long lives and documented their discoveries in detail inpapers, autobiographies, and interviews. These accounts offervaluable insights into the discovery process that can help guidenovel discoveries during periods of profound conceptual difficulties and confusion, such as the current state of artificial intelligence.

Therefore, I will quote the original writings of the main protagonists wherever appropriate, as I am convinced that theirown expressions lend authenticity and clarity to the very murkyprocesses behind great conceptual discoveries. My hope is togive the reader a sense of what is involved in achieving majorconceptual breakthroughs. As Max Planck said: In the history ofscience, a new concept never springs up in its complete and finalform, as in the ancient Greek myth, Pallas Athene sprang upfrom the head of Zeus. Heisenberg further elaborated: 5"The history of physics is not only a sequence of experimental discoveries and observations, followed bytheir mathematical description; it is also a history of concepts. For an understanding of the phenomena, the firstcondition is the introduction of adequate concepts. Only with the help of correct concepts can we really knowwhat has been observed. When we enter a new field, veryoften, new concepts are needed. As a rule, new concepts comeup in a rather unclear and undeveloped form. Later, they are modified, sometimes they are almost completely abandoned and are replaced by some better concepts, which then, finally, are clear and well-defined."

In the remainder of this Commentary, I provide a historical perspective that highlights key breakthroughs. This perspective is meant for those unfamiliar with quantum mechanics or itshistorical development. It is not aimed at experts. The objective of this paper is not to teach readers quantum mechanics but onlyto expose them to the central ideas, their historical evolution, and the conceptual struggles, with a moderate amount ofmathematics to illustrate these points. Given the scope of this perspective and its constraints, I will not discuss themathematical details, referring the readers to more comprehensive sources. 1,6–10 Furthermore, this is a personal perspective that reflects what I consider important and interesting developments. However, I believe that most quantum experts agree with the observations made in this paper.

2. TWO CLOUDS IN THE HORIZON: THE "1900-MOMENT"

At the dawn of the 20th century, on April 27, 1900, Lord Kelvin delivered an important lecture at the Royal Institution inLondon,11 summarizing the status of physics with the title "Nineteenth-Century Clouds Over the Dynamical Theory of Heat and Light." The "clouds" that bothered him were the twotroublesome experiments that did not agree with the theoretical predictions: (i) the null result of the Michelson-Morley experiment, which could not detect the motion of the Earththrough ether, and (ii) the ultraviolet catastrophe of blackbody radiation. Lord Kelvin correctly recognized the gravity of the situation and appreciated the profound uncertainty in the fundamentals of classical physics. As we know, these two "clouds" revolutionized physics, indeed all science, over the following three decades. 12

The first "cloud" led to the birth of the theory of relativity, completelyupending our understanding of space, time, gravity, and thecosmos itself. The second gave us quantum mechanics, openingthe secret door to an almost "magical" realm that we did not evenknow existed all around us all of the time. In fact, quantum theory was born soon in the same year, 1900, when Max Planckpresented his quantum hypothesis at a meeting of the GermanPhysical Society on December 14th, initiating the dispersal of the second cloud.13

This scientific drama unfolded like a well-written suspense thriller full of plot twists, turns, and surprising conceptual leaps, except that it was written in the language of mathematics, namely, linear algebra, differential equations, and probabilitytheory, echoing Galileo's declaration: 14"Philosophy is written in this grand book, I mean theuniverse, which stands continually open to our gaze, but it cannot be understood unless one first learns to comprehend the language in which it is written. It is written in the language of mathematics."In the annals of history, certain periods stand out as inflection points, times when scientific, technological, or social changes drastically altered the trajectory of our civilization. That momentin 1900, when Kelvin announced that all was not well in physics, was such a tipping point. The innovations that followed, both theory and in practical applications, continue to transform oursocieties and economies profoundly. There is no other thirtyyear period in history where our understanding of the universewas so dramatically upended as it was during 1900–1930.

3.ACT I: THE BIRTH OF QUANTUM THEORY (1900-1913)

Between 1900 and 1930, physicists were compelled to abandon classical mechanics in favor of quantum mechanics because the former could not predict or explain the atomic structure, spectrallines, and dual nature of matter and radiation as both waves and particles. This drama of frenetic intellectual activity occurred infour surprising breakthroughs. In the following sections, Iprovide an overview of these key advances.

3.1. Max Planck and Quantum Theory (1900). As noted, the roots of quantum mechanics can be traced to the "cloud"thatLord Kelvin worried about in the context of blackbody radiation. A blackbody is an idealized object that absorbs and emitselectromagnetic radiation at all frequencies. Classical physicspredicted the intensity of this radiation using the Rayleigh-Jeanslaw:

$$I(\lambda, T) = \frac{2ck_BT}{\lambda^4}$$
(1)

where $I(\lambda, T)$ is the radiation intensity, c is the speed of light, kB is Boltzmann's constant, T is the temperature, and λ is thewavelength of the radiation. This equation worked well at longwavelengths but diverged to infinity at short wavelengths, as shown in Figure 1, known as the ultraviolet catastrophe, a termcoined by Paul Ehrenfest in 1911.

Max Planck (Figure 2) got interested in this problem and, after a six-year struggle, introduced a revolutionary hypothesis:energy is not emitted continuously but in discrete Planck accomplished this in two critical steps, presented at the

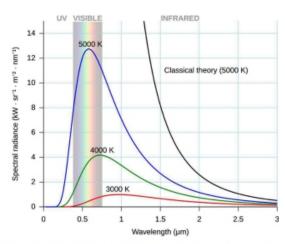


Figure 1. Ultraviolet catastrophe in classical Rayleigh-Jeans law. Reproduced with permission from ref 15. Copyright 2010 Darth Kule.

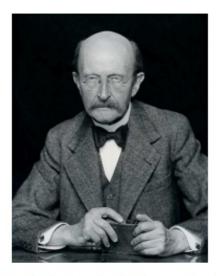


Figure 2. Max Planck. Reproduced with permission from ref 18. Copyright 1938 Hugo Erfurth.

German Physical Society meetings: (i) Discovering the correct radiation formula (on October 19, 1900) and (ii) Providing its conceptual justification via the quantum hypothesis (on December 14, 1900). Let us hear about the first step from Planck himself from hisscientific autobiography: 17"In fact, my previous studies of the Second Law of Thermodynamics came to stand me in good stead now, for at the very outset I hit upon the idea of correlating notthe temperature but the entropy of the oscillator with its energy. It was an odd jest of fate that a circumstance whichon former occasions I had found unpleasant, namely, thelack of interest of my colleagues in the direction taken by myinvestigations, now turned out to be an outright boon. Whilea host of outstanding physicists worked on the problem of spectral energy distribution, from both the experimental andtheoretical aspects, every one of them directed his efforts solely toward exhibiting the dependence of the intensity of radiation on the temperature On the other hand, I suspected that the fundamental connection lies in the dependence of entropy upon energy. As the significance of the concept of entropy had not yet come tobe fully appreciated, nobody paid any attention to themethod adopted by me, and I could work out mycalculations completely at my leisure, with absolutethoroughness, without fear of interference or competition ... In this way, a new radiation formula was obtained, and Isubmitted it for examination to the Berlin Physical Society, at the meeting on October 19, 1900." Although Rudolf Clausius introduced the concept of entropyin 1864, it remained undervalued by the scientific community, surprisingly, for nearly three decades. This highlights the significant amount of time required for revolutionary to gain widespread acceptance. Reflecting on this, Planck later remarked rather sardonically:17"A new scientific truth does not triumph by convincing itsopponents and making them see the light, but rather because its opponents eventually die out, and a new generation growsup that is familiar with it." Planck's new radiation law is given by

$$I(\lambda, T) = \frac{2hc^2}{\lambda^5} \frac{1}{e^{hc/\lambda k_B T} - 1}$$
(2)

(h is Planck's constant) correctly described experimental data at all wavelengths and resolved the ultraviolet catastrophe. However, Planck was unsatisfied with the clever guessworkthat led to its discovery. He wanted to know its significance, and so he proceeded with the second step:17" But even if the absolutely precise validity of the radiation formula is taken for granted, so long as it had merely the standing of a law disclosed by a lucky intuition, it could not be expected to possess more than a formal significance. For this reason, on the very day that I formulated this law, Ibegan to devote myself to the task of investing it with a truephysical meaning. This quest automatically led me to study the interrelation

of entropy and probability in other words, to pursue the line of thought inaugurated by Boltzmann. Since the entropy S is an additive magnitude but the probability W is a multiplicative one, I simply postulated that $S = k \cdot \log W$, where k is a universal constant; and I investigated whether the formula for W, which is obtained when S is replaced by its value corresponding to the aboveradiation law, could be interpreted as a measure of probability ... It is, understandably, often called Boltzmann's constant. However, this calls for the comment that Boltzmann never introduced this constant, nor, to the best of my knowledge, did he ever think of investigating its numerical value." Interestingly, the famous equation S = k log W, which is inscribed on Boltzmann's tomb in Vienna, was not stated in this form by Boltzmann. It was Planck who expressed the Boltzmannresult in this now familiar form.16Applying Boltzmann's reasoning about entropy from his 1877paper19 to blackbody radiation, Planck was led to the concept of discrete packets of energy, which he termed quanta. 13,20 He was, however, uncomfortable with this idea as he was aware that hewas violating the continuity principle, 21 a fundamental principle that dates back to Leibnitz, who famously said:22 "Natura nonfacit saltus" (Latin for "nature does not make jumps"). This principle also serves as the foundation for differential and integral calculus. Planck alerts us to this crucial feature of his theory in hisDecember 14, 1900, paper:13"If E [the total energy] is considered to be a continuous divisible quantity, this distribution is possible in infinitely many ways. We consider, however this is the mostessential point of the whole calculation E to be composed a well defined number of equal parts [of magnitude ϵ] and use thereto the constant of nature h = $6.55 \times 10-27$ ergsec [setting $\epsilon = hv$]." About his break with classical physics tradition and his embracing of Boltzmann's "atomistic" ideas, which he had beencritical of for many years, Planck would later recall: "Briefly summarized, what I did can be described as simplyan act of desperation. By nature, I am peacefully inclined and reject all doubtful adventures. But by then I had been wrestling unsuccessfully for six years (since 1894) with the problem of equilibrium between radiation and matter, and Iknew that this problem was of fundamental importance tophysics; I also knew the formula that expresses the energy distribution in normal spectra. A theoretical interpretation therefore had to be found at any cost, no matter how high. It was clear to me that classical physics could offer nosolution to this problem and would have meant that all ofthe energy would eventually transfer from matter intoradiation. In order to prevent this, a new constant is required to ensure that energy does not disintegrate. But the only way to recognize how this can be done is to start from adefinite point of view. This approach was opened to me bymaintaining the laws of thermodynamics. The two laws, itseems to me, must be upheld under all circumstances. Forthe rest, I was ready to sacrifice every one of my previous convictions about physical laws." Although historians continue to debate how much Planckrealized the significance of his quantum hypothesis, 13, 20, 21 Erwin Planck later recalled what his father told him soon after hisdiscovery:20"Either what I have found out now is complete nonsense orit might be one of the greatest discoveries in physics sinceNewton."

Planck proposed that the energy of a radiation mode isquantized and proportional to its frequency:

$$E = nh\nu$$
 $n = 1, 2, 3, ...$ (3)

where h is Planck's constant and v is the frequency of the radiation. This assumption prevented infinite energy at shortwavelengths and correctly described blackbody radiation. Usingenergy quantization, Planck derived the formula presented above, now expressed in terms of frequency v:

$$I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/k_BT} - 1}$$
(4)

Max Planck was awarded the 1918 Physics Nobel Prize for his discovery. Although Planck gave birth to quantum theory, heremained a reluctant revolutionary for a long time, viewingquantization only as a "mathematical trick" rather than as afundamental law of nature. After all, he was very much a part of the old guard and therefore was hesitant to abandon the "sacred" principle of continuity. It was only much later that he camearound to accepting the new reality.

3.2. Albert Einstein and the Photoelectric Effect (1905).

On the other hand, the Young Turk who followed Planck next in this exciting drama was a rebellious iconoclast who was rearing to upend the very foundations of physics, notjust radiation theory. Enter Albert Einstein (Figure 3), a 26-year



Figure 3. Albert Einstein. Reproduced with permission from ref 23. Copyright 1905 Lucien Chavan.

old unknown clerk at the Swiss Patent Office in Bern. While Planck saw quantization as just a "mathematical trick", not afeature of physical nature, Einstein took it more seriously as afundamental property of nature. He was open to such radical rethinking as he was, at the same time, busy overthrowing the Newtonian concepts of space and time in his new theory of special relativity, which dispersed the first "cloud" that LordKelvin worried about.

In his Annus Mirabilis, 1905, Einstein extended Planck's idea and proposed that light itself consists of quantized particlescalled photons, each carrying energy E = hv, to solve a puzzlingresult in the photoelectric effect, where light incident on a metalsurface ejects electrons. The classical wave theory of lightpredicted that increasing the light intensity should increase theelectron energy. However, experiments showed that noelectrons are emitted below a threshold frequency, regardlessof intensity. They further showed that the electron energydepends on the frequency, not the intensity. The higher lightintensity increased the number of emitted electrons, but nottheir individual energy. Einstein's theory treats the photoelectric effect as a one-toone interaction between a photon and an electron. The energybalance equation is

$$h\nu = W + E_k$$
 (5)

where hv is the energy of the incoming photon, W (or work function) is the minimum energy required to free an electronfrom the metal, and Ek is the maximum kinetic energy of theemitted electron. This equation explained all of the experimental results. Reflecting on his ground breaking papers from 1905, which included his first two on relativity, he regarded only the light-quanta paper as genuinely revolutionary. Einstein wasawarded the Nobel Prize in Physics in 1921 for this discovery. It is interesting to note that the prize was not for his work on relativity. From the perspective of the evolution of quantum theory, Einstein's theory confirmed the particle nature of light, supporting the idea that light exhibits both

wave and particle properties, i.e., wave-particle duality. This paved the way for theacts that followed next.

3.3.Niels Bohr and theHydrogen Atom (1913). In 1897, J. J. Thomson at Cambridge discovered the electron in cathoderay tube experiments. Ernest Rutherford, who had trained underThomson, showed in his gold foil experiment at Manchester in1911 that the atom is mostly an empty space with a tiny, dense, positively charged nucleus with orbiting electrons. However, classical electrodynamics predicted that electrons should spiralinto the nucleus due to radiation loss. Furthermore, the atomicspectra of hydrogen showed discrete spectral lines, contradicting classical physics, which predicted a continuous spectrum. In 1911, Niels Bohr (Figure 4) arrived in England to study atomic structure under Thomson first and with Rutherford later



Figure 4. Niels Bohr. Reproduced with permission from ref 24. Copyright 1922 AB Lagrelius & Westphal.

in Manchester. In 1913, in a series of three papers, Bohr proposed an atomic model that resolved the contradictions. While Einstein extended Planck's quantum hypothesis tophotons, Bohr further extended it to electrons by introducing quantized orbits for electrons. Bohr's atomic model introduced three key quantum postulates:25

(i) Electrons move in fixed circular orbits around the nucleus, where their angular momentum is quantized:

$$L = n\hbar = n\frac{h}{2\pi}, \quad n = 1, 2, 3, ...$$
 (6)

where L is the electron's angular momentum, $\hbar = h/2\pi$ is the reduced Planck constant, and n is the principal quantum numberspecifying the allowed orbits. This assumption prevented electrons from spiraling into the nucleus, ensuring atomic stability.

(ii) The total energy of an electron in orbit is also quantized and given by

$$E_n = -\frac{13.6\text{eV}}{n^2} \tag{7}$$

where En is the energy of an electron in orbit n, and -13.6 eV is the ground-state energy of hydrogen (energy levels are negative, meaning that electrons are bound to the nucleus). This quantization explains why the atoms do not radiate continuously.

(iii) Electrons can transition between orbits, i.e., performquantum jumps, by absorbing or emitting a

photon of energy:

$$E_{\text{photon}} = h\nu = E_i - E_f \qquad (8)$$

where v is the frequency of emitted/absorbed light and Ei and Ef are the initial and final energy levels. This correctly explainedhydrogen's spectral lines, known as the Balmer series, given by

$$\frac{1}{\lambda} = R_H \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right) \tag{9}$$

where RH is the Rydberg constant.

Thus, the Bohr model successfully explained the atomic stability and correctly predicted hydrogen spectral lines.26 Mostimportantly, Bohr had conceptually generalized Planck's "mathematical trick" and made quantization a fundamental feature of nature. Bohr was awarded the Nobel Prize in Physicsin 1922. With this, the first act ends, and the stage is set for even more surprising twists and turns.

4. ACT II: DE BROGLIE AND WAVE-PARTICLE DUALITY (1923-1924)

About ten years after Bohr, the next crucial conceptual breakthrough came in the form of further generalization of thewave-particle duality of light. In 1923, Prince Louis de Broglie(Figure 5) introduced the shocking concept of matter waves,



Figure 5. Louis de Broglie. Reproduced with permission from ref 29. Copyright 1929 University of Maryland

which impressed Einstein so much that he remarked:27 "He has lifted a corner of the great veil." Put simply, de Broglie askedhimself: If light, which was thought of as a wave, can exhibit particle-like behavior (as photons), why cannot particles likeelectrons exhibit wave-like behavior? He defended this idea inhis Ph.D. thesis on 25 November 1924, in Paris.28 In this thesis,he proposed wave-particle duality for matter, suggesting that "matter waves" obey the equation

$$\lambda = \frac{h}{p}$$
(10)

where λ is the de Broglie wavelength and p is momentum. In 1927, Davisson and Germer, and independently Thomson and Reid, confirmed this idea in electron diffraction experiments. deBroglie received his Nobel Prize in Physics in 1929, a mere fiveyears after his Ph.D. defense. Davisson and Thomson received theirs in 1937.

One question physicists and historians have puzzled over for many years is why de Broglie, who discovered matter waves, didnot proceed to discover Schrödinger's wave equation. Althoughwe cannot be certain, experts have identified several reasons afterconducting careful studies. Here, I quote Olivier Darrigol:27"A first element of the answer is that, notwithstanding withhis grand analogy between dynamics and optics, he (deBroglie) was shy in adventuring beyond the approximation of geometrical optics. He focused on retrieving results of thereceived quantum theory, such as the Bohr–Sommerfeldconditions, and he underplayed the more disturbing consequences of his concept of matter waves.

Another possible obstacle to his developing a wave theory of matter was his conviction that both light and matter had adual nature, implying the synchronous motion of waves and particles. This duality focused on the interplay betweenwaves and particles rather than on the search for a newwave equation. Third and most importantly, de Broglie believed that the analogy between light and matter implied the electromagnetic nature of his matter waves. Consequently, he also believed that matter waves obeyed the d'Alembertian equation of electromagnetism. Direct evidence of this conviction is found in a note of 1925 in which he describes the intrinsic oscillation of an electron in its rest frame as the stationary superposition of the retarded and advanced solutions of the d'Alembertian equation. The same heuristic principle, the analogy between matterand light, led de Broglie to the matter waves and prevented him from seeking a specific equation for these waves!"

Therein lies a very important lesson in the use of analogies to discover new conceptual breakthroughs. One should not take it too literally or expect an exact analogy of the new phenomenonin every detail. Although de Broglie was correct in reasoning that the wave-particle duality of light implied a similar duality forelectrons (matter, in general), he took this analogy too far to reason that matter waves would also be electromagnetic innature. This is where the analogy broke down. Fortunately, Schrödinger did not make this mistake!

5. ACT III: THE BIRTH OF QUANTUM MECHANICS (1925–1927)

Finally, we arrive at the main event, the birth of quantum mechanics. The key characters are Werner Heisenberg, ErwinSchrödinger, Max Born, Paul Dirac, and Wolfgang Pauli. Even adecade after the Bohr atom, atomic phenomena have remainedlargely unexplained, with many disturbing fundamental questions. There was no coherent mathematical theory yet, only a collection of seemingly ad hoc rules of quantum behavior. The transition from classical mechanics to quantum mechanics remained an elusive goal before 1925.

5.1. Heisenberg and Born: Matrix Mechanics. The first major breakthrough in resolving this impasse was initiated by23-year-old Heisenberg (Figure 6) in his historic 1925 papernoted earlier,3 marking the birth of quantum mechanics. Heisenberg's innovative idea, guided by Bohr's CorrespondencePrinciple, was to retain classical mechanics equations but replacethe classical position coordinate with a quantum-theoretical quantity. The new position quantity contains information about the measurable line spectrum of an atom rather than theunobservable orbital of the electron. He devised a special kinematical rule for multiplying position quantities. Mehra gives a vivid description of this momentous discovery:



Figure 6. Werner Heisenberg. Reproduced with permission from ref 31 Copyright 1933 German Federal Archives.

"With the coming of spring in 1925, Heisenberg had developed a case of severe hay fever, which would just notleave him, and he decided to take a week or ten days off inJune 1925 at the rocky island of Helgoland in the NorthSea. At Helgoland, not only did he cure his hay fever butwiped the nose clean of the chronic colds of erstwhileproblems of atomic mechanics ...At Helgoland, Heisenberg divided his time in taking longwalks, reading Goethe's West-Ostlicher Divan, and seekingto give his vague ideas on quantum mechanics a moredefinite shape. There he solved two problems ...The example of the anharmonic oscillator showed him that a dynamical problem in quantum theory could be solved with the help of his scheme.

As he (Heisenberg) recalled:30 'It was almost three o'clock in the morning before the final result of my computations laybefore me. The energy principle had held for all of the terms, and I could no longer doubt the mathematical consistency and coherence of the kind of quantum mechanics to whichmy calculations pointed. At first, I was deeply alarmed. had the feeling that, through the surface of atomic phenomena, I was looking at a strangely beautiful interiorand felt almost giddy at the thought that I now had to probe this wealth of mathematical structures nature had sogenerously spread out before me. I was far too excited to sleep, and so, as a new day dawned, I made a trip to the southern tip of the island, where I had been longing to climba rock jutting out into the sea. I now did so without too much trouble, and waited for the sun to rise." After he returned from Helgoland, Heisenberg gave his paperto Max Born (Figure 7) in early July for his opinion. Heisenbergwas working as Born's research assistant at the University of Gottingen at that time. Born had been keenly aware of the difficulties in quantum theory for some time as he wrote:32 "It becomes increasingly probable that not only new assumptions will be needed in the sense of physical hypotheses, but that the entire system of concepts of physics must be rebuilt from the ground up." So, when he saw Heisenberg's new mathematical



Figure 7. Max Born. Reproduced with permission from ref 33. Copyright 1954 German Federal Archives.

formulation of kinematics of quantum systems, Born immediately recognized its importance, as he recalls:2 "I began to ponder about his symbolic multiplication andwas soon involved in it. I thought the whole day and couldhardly sleep at night ... In the morning I suddenly saw thelight: Heisenberg's symbolic multiplication was nothing butthe matrix calculus, well-known to me since my student daysfrom the lectures of Rosanes in Breslau." A few days later, on July 19, 1925, Born traveled from Göttingen to Hanover to attend a meeting of the German Physical Society, where he informed Wolfgang Pauli about thematrices. Pauli was critical:2"Yes, I know that you are fond of a tedious and complicated formalism. You are only going to spoil Heisenberg's physicalideas by your futile mathematics." To a modern physicist, it is astonishing that Heisenberg didnot know about matrices when he made his great discovery, as headmits:5"At that time I must confess I did not know what a matrixwas and did not know the rules of matrix multiplication." As Fedak and Prentis describe, 34 it was Born who recognized that the next step was to formalize Heisenberg's theory using the language of matrices, which he did with his student PascualJordan35 after Pauli turned him down.2 This was followed by another paper by Born, Heisenberg, and Jordan. 36 It was also Born who coined the name Quantum Mechanics for the newfield.34,37 Born expressed Heisenberg's results in a more elegantform using the matrix notation. If Q and P are the position and momentum matrices, they satisfy

$$[P, Q] = PQ - QP = (h/2\pi i)I$$
 (11)

where I is the identity matrix, and the quantity [P,Q] is known as the commutator. It is important to note that using matrices is notjust a matter of mathematical elegance. What Heisenberg haddiscovered inadvertently was one of the fundamental aspects of quantum reality: its dynamic variables are represented by operators (and hence matrices), unlike classical variables, which are represented by scalars. This critical feature was independently recognized by Paul Dirac around the same time 38 (more on this below). These papers introduced a novelapproach to atomic Hamiltonian mechanics using non commutative quantum methods. This marked the beginning of a new phase in theoretical physics, characterized by the use of Hermitian matrices, commutators, and eigenvalue problems askey mathematical tools in atomic theory.

This noncommutativity of position and momentum matrices led to a major breakthrough two years later, in 1927, whileHeisenberg was visiting the Niels Bohr Institute in Copenhagen.He describes what happened one late evening as he took a strollthrough Faelledparken, the lovely park behind the

institute:30

"It must have been one evening after midnight when Isuddenly remembered my conversation with Einstein and particularly his statement, 'It is the theory which decides what we can observe.' I was immediately convinced that thekey to the gate that had been closed for so long must be sought right here. I decided to go on a nocturnal walkthrough Faelled Park and to think further about the matter. We had always said so glibly that the path of the electron in the cloud chamber could be observed. But perhaps what were ally observed was something much less. Perhaps we merely a series of discrete and ill-defined spots through which the electron had passed. In fact, all we do see in the cloudchamber are individual water droplets, which must certainly be much larger than the electron. The right question should therefore be: Can quantum mechanics represent the fact that an electron finds itself approximately in a given place and that it moves approximately with a given velocity, and can we make these approximations so close that they do not cause experimental difficulties?

A brief calculation after my return to the Institute showed that one could indeed represent such situations mathematically and that the approximations are governed by whatwould later be called the uncertainty principle of quantummechanics: the product of the uncertainties in the measuredvalues of the position and momentum (i.e., the product ofmass and velocity) cannot be smaller than Planck'sconstant. This formulation, I felt, established the muchneeded bridge between cloud chamber observations and themathematics of quantum mechanics. True, it had still to be proved that any experiment whatsoever was bound to set upsituations satisfying the uncertainty principle, but this struckme as plausible a priori since the processes involved in the experiment or the observation had necessarily to satisfy the laws of quantum mechanics. On this presupposition, experiments are unlikely to produce situations that do accord with quantum mechanics. 'It is the theory which

decides what we can observe.' I resolved to prove this bycalculations based on simple experiments during the nextfew days."The uncertainty principle states that there is an intrinsic limit how precisely we can simultaneously measure the position qand the momentum p of a particle. Heisenberg derived the following inequality:

$$\Delta q \cdot \Delta p \ge \frac{\hbar}{2}$$
(12)

where Δq is the standard deviation of position and Δp is the standard deviation of momentum. If we try to measure aparticle's position very precisely (Δq small), the uncertainty immomentum Δp increases. Conversely, if we measure themomentum precisely, the uncertainty in the position grows. This principle is not due to measurement errors but rather aninherent property of quantum systems.

Given this history of the uncertainty principle and its close association with the Niels Bohr Institute, I found it so fitting, in alighter vein, to see this cartoon (Figure 8) displayed on a door of the Institute during my visit in August of 2022.



Figure 8. Cartoon on a wall of the Niels Bohr Institute. Photo by the author in 2022. Artist unknown.

There are serious implications captured by this fundamental property of nature. (i) Observer's interference: the very act measurement disturbs the system. (ii) Wave-particle duality: position and

momentum cannot be simultaneously well-defined.(iii) Limits of classical concepts: the classical idea of a trajectorydoes not hold in the quantum realm.

Heisenberg would later speak in sheer awe of the startling simplicity and beauty of the new theory:30"If nature leads us to mathematical forms of great simplicity and beauty by forms, I am referring to coherent systems of hypotheses, axioms, etc. • to forms that no one has previously encountered, we cannot help thinking that they are 'true,' that they reveal a genuine feature of nature ... Youmust have felt this too: the almost frightening simplicity andwholeness of the relationships which nature suddenlyspreads out before us and for which none of us was in theleast prepared."Heisenberg was awarded the 1932 Nobel Prize in Physics. Given Born and Jordan's pivotal role in the discovery ofquantum mechanics, it is natural to wonder why they were leftout. In 1933, Heisenberg wrote Born saying:39"The fact that I am to receive the Nobel Prize alone, forwork done in Göttingen in collaboration you, Jordan, and I this fact depresses me, and I hardly know what to write to you. I am, of course, glad that our common now appreciated and I enjoy the recollection of the beautiful time of collaboration. I also believe that all good physicistsknow how great was your and Jordan's contribution to the structure of quantum mechanics and this remains unchanged by a wrong decision from outside. Yet I myselfcan do nothing but thank you again for all the finecollaboration and feel a little ashamed."Fortunately, Born was awarded the Nobel Prize in Physics in1954 for his fundamental research in quantum mechanics, especially for his statistical interpretation of the wave function (as discussed below). Engraved on Max Born's tombstone in Göttingen is a one-line epitaph: pq - qp = $h/2\pi i$.

5.2. Schrödinger's Wave Mechanics (1926). 1925 was already an amazing year, but the quantum mechanics revolutionwas not yet finished for the year. Following a line of attack that is different from the matrix mechanics formalism, ErwinSchrödinger (Figure 9) was developing something veryinteresting. Inspired by de Broglie's matter waves, he introducedwave mechanics, and the fundamental equation governing quantum evolution, the Schrödinger equation:



Figure 9. Erwin Schrödinger. Reproduced with permission from ref 40. Copyright 1930 Nobel Foundation.

$$i\hbar \frac{\partial}{\partial t} \psi = \hat{H} \psi$$
(13)

where \hat{H} is the Hamiltonian operator and ψ is the wave function. The time-independent version,

$$-\frac{\hbar^2}{2m}\nabla^2\psi + V\psi = E\psi \qquad (14)$$

explains energy quantization and atomic structure. Schrödinger showed that wave mechanics is mathematically equivalent tomatrix mechanics.

Just as he proposed the matrix formalism to clarify Heisenberg's quantum mechanics, Max Born once again steppedup and clarified the meaning of the wave function in wavemechanics in 1926.41 Born interpreted the wave function $\psi(x, t)$ as a probability amplitude. The probability of finding a particle at position x is given by:

$$P(x) = |\psi(x)|^2 \tag{15}$$

This marked a fundamental conceptual shift from a deterministic perspective of the universe in classical mechanics a probabilistic view of quantum mechanics. It is indeed quiteremarkable that such a fundamental interpretation that completely revolutionized our view of the universe wasmentioned in a mere footnote of Born's 1926 paper.41 In fact, there is a fascinating backstory to this. In a paper written on theoccasion of the birth centenary of Born in 1982, Abraham Paisobserved:42

"Then, Born declares: ' ϕ mn (i.e. the wavefunction, $\psi(x)$) determines the probability for the scattering of the electronfrom the z-direction into the direction $[\theta, \phi]$.' At best, this statement is vague. Born added a footnote inproof to his evidently hastily written paper: 'A more preciseconsideration shows that the probability is proportional tothe square of ϕ mn.' He should have said 'absolute square.' But he clearly had got the point, and so the correct expression for the transition probability concept enteredphysics via a footnote.

I shall return shortly to the significant fact that originally associated probability with ϕ mn rather than with $|\phi$ mn|2. As I learned from recent private discussions, Dirachad the very same idea at that time. So did Wigner, whotold me that some sort of probability interpretation was thenon the minds of several people, and that he, too, had thoughtof identifying ϕ mn or $|\phi$ mn| with a probability. When Born'spaper came out and $|\phi$ mn|2 turned out to be the relevant quantity, 'I was at first taken aback but soon realized thatBorn was right,' Wigner said."It is absolutely incredible and deeply instructive that such afundamental feature of quantum mechanics, namely, its probabilistic nature, was initially guessed wrong even by giantslike Born, Dirac, and Wigner, and was subsequently corrected ina footnote only during the proof stage of the manuscript. Again, this teaches us valuable lessons about the nature of the discoveryprocess, particularly fundamental concepts.

Over the years, many have wondered why Schrödinger, of alltheoretical physicists, took up de Broglie's ideas and developed them into wave mechanics. 43 We briefly saw above why deBroglie himself did not do it. Raman and Forman provide an interesting account 43 that de Broglie was not taken seriously by the quantum establishment:

"Thus in Copenhagen and in Gottingen, where atomicphysics was pursued in the Copenhagen spirit, de Brogliewould certainly have had the reputation of a renegade, if notexactly a crank, who stuck obstinately to his own illconceived theories ... Thus among the central Europeanphysicists deeply involved in the problems of theoreticalspectroscopy, and this was indeed the great majority of thoseseriously concerned with the quantum theory, de Brogliemust have had a very bad reputation."On the other hand, Schrödinger had no such biases against deBroglie and so took his work seriously. There is a well-knownanecdote due to Dirac44 that the first wave equation Schrödingerguessed later became known as the relativistic Klein-Gordonequation. When this equation, applied to the hydrogen atom, didnot yield the familiar results, Schrödinger abandoned thisequation, searched again for a better candidate, and discoveredthe famous Schrödinger equation. Felix Bloch, the 1952 Nobel laureate in Physics, who was astudent at ETH-Zurich at that time, provides additional details45on the events when Schrödinger participated in their physicscolloquium run by Peter Debye (Nobel Prize in Chemistry,1936). Bloch recalls Schrödinger's seminar in early 1925:

"Once at the end of a colloquium I heard Debye saying something like: 'Schrödinger, you are not working right nowon very important problems anyway. Why don't you tell ussome time about that thesis of de Broglie, which seems tohave attracted some attention.'So, in one of the next colloquia, Schrödinger gave a beautifully clear account of how de Broglie associated awave with a particle and how he could obtain thequantization rules of Niels Bohr and Sommerfeld bydemanding that an integer number of waves should be fittedalong a stationary orbit. When he had finished, Debyecasually remarked that he thought that this way of talkingwas rather childish. As a student of Sommerfeld he hadlearned that, to deal properly with waves, one had to have awave equation ... Just a few weeks later he (Schrödinger)gave another talk in the colloquium which he started bysaying: 'My colleague Debye suggested that one should have a wave equation; well, I have found one!" Prompted by Debye, Schrödinger discovered his equation in about three months, between November 1925 and

January 1926, and published a series of four papers on wavemechanics entitled Quantization as an eigenvalue Problem.46–48It is understandable that the members of the "CopenhagenEstablishment" did not discover the wave equation, as they did not take de Broglie seriously. But I have often wondered why Einstein or Debye did not discover the wave equationthemselves. I believe that while Einstein understood theimportance of de Broglie's matter wave concept, he was toopreoccupied with his search for the unified field theory, which heworked on for the rest of his life. As for Debye, it appears that hehad some regrets, as narrated again by Bloch:45"Many years later, I reminded Debye of his remark aboutthe wave equation; interestingly enough he claimed that hehad forgotten about it, and I am not quite sure whether thiswas not the subconscious suppression of his regret that hehad not done it himself. In any event, he turned to me with abroad smile and said: 'Well, wasn't I right?'""

Initially, Heisenberg's matrix mechanics and Schrödinger's wave mechanics appeared to be very different from each other, and an acrimonious debate ensued over which one was correct. In a footnote to a 1926 paper, Schrödinger wrote: "I wasdiscouraged, if not repelled, by what appeared to me rather difficult method of transcendental algebra, defying any visualization." Meanwhile, Heisenberg complained to Pauli: "Themore I think about the physical part of Schrödinger theory, themore detestable I find it." Fortunately, the debate was resolved in 1926. Schrödinger, along with Carl Eckert, working independently, demonstrated that the two new mechanics, although superficially very different, were mathematically equivalent to each other. Schrödinger was awarded the Nobel Prize in Physics in 1933, which he shared with Paul Dirac, discussed next, for their contributions to quantum mechanics. 5.3. Commutator and the Poisson Brackets: Dirac's Discovery (1928–1930). Right before the paper by Born, Heisenberg, and Jordan was published in January 1926, another paper outlining the whole framework of quantum mechanics was published in the Proceedings of the Royal Society by Paul Dirac (Figure 10), then a research student of R. H. Fowler's in Cambridge. Reflecting on Heisenberg's paper, Dirac recalled: 49



Figure 10. Paul Dirac. Reproduced with permission from ref 50. Copyright 1933 Nobel Foundation.

"During a long walk on a Sunday it occurred to me that the commutator might be the analogue of the Poisson bracket, but I did not know very well then what a Poisson bracketwas. I had just read a bit about it and forgotten most ofwhat I had read. I wanted to check up on this idea, but Icould not do so because I did not have any book at homethat gave Poisson brackets, and all the libraries were closed. So I had just to wait impatiently until Monday morningwhen the libraries were open to check on what Poissonbracket really was. Then I found that they would fit, but Ihad one impatient night of waiting." By recognizing the link between these two brackets, Diraceffectively clarified the connection between Heisenberg's variables and classical variables, giving the formulation a moreclassical appearance. Meanwhile, it neatly highlighted the precisepoint where the reformulation diverged from the classical theory. Dirac was one of the most brilliant theoretical physicists of thetwentieth century, making profound contributions to quantummechanics, quantum field theory, and relativistic quantummechanics. His work introduced the Dirac equation, predicted the existence of antimatter, and laid the mathematical foundation for quantum electrodynamics (QED). Dirac sharedhis Nobel in 1933 with Schrödinger.

5.4. Pauli Exclusion Principle (1925). As I wrap up this period of frenetic activity, I would be remiss if I did not mention the contributions of Wolfgang Pauli (Figure 11), particularly his exclusion principle. Pauli made fundamental contributions to quantum mechanics and quantum field theory, significantly shaping modern physics. His most famous work includes the Pauli exclusion principle, his contributions to spin theory, the theory of quantum electrodynamics (QED), and the prediction of the neutrino. In 1925, Pauli formulated the exclusion principle, stating that no two identical fermions can occupy the same quantum states imultaneously. Mathematically, this means that for a system of two electrons, the wave function Ψ must be antisymmetric underparticle exchange: $\Psi(1,2) = -\Psi(2,1)$. This ensures that if two electrons were in the same quantum state, then the wavefunction would be zero, prohibiting such configurations.



Figure 11. Wolfgang Pauli. Reproduced with permission from ref 51. Copyright 1945 Nobel Foundation.

The Pauli exclusion principle explains: (i) electron shell structure of atoms, (ii) periodic table organization and whydifferent elements have distinct chemical properties, and (iii)stability of matter, as it prevents electrons from collapsing into the lowest energy state. For his contributions to the development of quantum mechanics, Pauli was awarded the Nobel Prizein Physics in 1945.

6. ACT IV: THE COPENHAGEN INTERPRETATION (1927-1930)

Starting with Heisenberg's matrix mechanics in 1925 and concluding with Dirac's relativistic quantum theory in 1930, in ashort span of five years, a coherent mathematical formalism of quantum mechanics emerged. However, its conceptualimplications seriously bothered several leading physicists, including those who contributed to its development, such as Einstein, Schrödinger, and others. Objecting to the probabilistic foundations of quantum mechanics, Einstein was perhaps themost vocal, famously saying:52 "God does not play dice with the universe." On quantum entanglement,52 he called it "spookyaction at a distance." Schrödinger devised the famous Schrödinger's cat paradox to highlight the interpretational issues of quantum mechanics.

Despite such objections, physicists converged around a set of principles advocated by Bohr and Heisenberg in 1927, known as the Copenhagen Interpretation, which has remained the most widely accepted view of quantum mechanics for a century. Thekey tenets of this view are: (i) Nature at the quantum level is intrinsically probabilistic, and the square of the wave function $|\psi(x,t)|$

2 gives the probability of finding a particle at (x, t). (ii) A quantum system exists in a superposition until measured, atwhich point it collapses into a definite state. (iii) The act ofmeasurement affects the system. (iv) Key quantities such asenergy, momentum, spin, etc. are quantized. There are some fundamental concerns with this interpretation of quantum mechanics, particularly with respect to the wavefunction collapse, which we shall not go into .53–56 The fact that the predictions of quantum mechanics have been fantastically accurate, as verified by countless experiments over the decades, although its conceptual foundations are somewhat murky, prompted N. David Mermin, the physics professor who taught me quantum mechanics at Cornell, to summarize the Copenhagen Interpretation as "Shut up and calculate!" This quote is often misattributed to Richard Feynman.

7. IMPACT OF QUANTUM MECHANICS IN CHEMICAL ENGINEERING

Although the objective of this paper is not on the application of quantum mechanics, I would like to briefly mention its profoundimpact on chemical engineering and materials science. 58,59 From reaction kinetics to materials design, quantum mechanicsprovides the fundamental principles that govern atomicinteractions, electronic structure, chemical bonding, computational chemistry, catalysis, nanotechnology, and quantum computing, among other areas. Quantum mechanics provides insights into (i) molecular interactions and reaction mechanisms, (ii) electronic structures governing chemical and material properties, and (iii) energy levels that define molecular and solid-state behaviors. Using such information, chemicalengineers optimize catalysts, polymers, drug molecules, andnanomaterials, improving efficiency and sustainability. For example, the Schrödinger equation is routinely used todetermine molecular structures and properties, such as bondlengths and angles, reaction energy barriers for kinetic analysis, and molecular orbitals and charge distributions. The DensityFunctional Theory is widely used to design catalysts, semiconductors, polymers, and nanomaterials. Quantum dots are yetanother application for designing nanoscale semiconductors with tunable electronic properties used in LED displays and photovoltaics. Quantum confinement is utilized, for example, in the design of graphene-based sensors and supercapacitors forenergy storage. As quantum technology advances, chemicalengineering and materials science will continue to leverage itsprinciples for sustainable industrial processes, advanced materials, and novel pharmaceuticals, driving innovation in the 21st century.

8. ISAIATA "1900-MOMENT"?

From its origins in abstract thought to its applications in materials science and quantum computing,

quantum mechanicsis a testament to the power of the human intellect to unlock nature's most closely guarded secrets. Quantum mechanicsrevolutionized physics by fundamentally altering our understanding of nature on the atomic scale. As Bohr remarked: "Ifquantum mechanics has not profoundly shocked you, youhaven't understood it yet." The key conceptual breakthroughs, summarized in Table 1, reveal an interesting finding. It appears that even the pioneersmissed the next conceptualstep. For example, Planck consideredhis quantum hypothesis merely a "mathematical trick," not afundamental law of nature, and, therefore, missed the connection with the photoelectric effect. Einstein understoodthis connection, but surprisingly, he did not realize its implications for other kinds of matter when he applied thehypothesis to photons. It was Bohr who connected it to electrons and their atomic orbitals, yet he, too, failed to grasp itsgenerality. de Broglie was the one who perceived the universalnature of the wave-particle duality. However, his excessivereliance on electromagnetic wave analogies prevented him from discovering the wave equation, a feat accomplished by Schrödinger. Again, Schrödinger did not quite understand conceptual significance of the wave function, which Born later

Table 1. Key Developments in Quantum Mechanics (1900–1930)

-,50,	
Year	Development
1900	Following Boltzmann's reasoning, Planck proposes his quantum hypothesis: Energy is quantized in discrete packets (quanta, $E=h\nu$).
1905	Einstein's photoelectric effect: Light behaves as particles (photons) with energy $(E=h\nu)$.
1913	Bohr's atomic model: Electrons exist in quantized orbits, explaining hydrogen spectra.
1924	de Broglie's wave-particle duality: Matter exhibits both wave-like and particle-like properties.
1925	Heisenberg's matrix mechanics: The first mathematical formulation of quantum mechanics.
1925	Pauli Exclusion Principle: No two identical fermions (e.g., electrons) can occupy the same quantum state simultaneously, explaining the structure of electron shells in atoms.
1926	Schrödinger's wave equation: Describes quantum states using wave functions.
1926	Born's probabilistic interpretation: The absolute square of the wave function represents probability amplitudes, introducing the statistical nature of quantum mechanics.
1927	Heisenberg's uncertainty principle: Position and momentum cannot be precisely known simultaneously.
1927	The Copenhagen Interpretation: Quantum mechanics is fundamentally probabilistic. The wave function collapses upon measurement, and complementarity dictates that quantum objects exhibit either particle or wave-like behavior depending on observation.
1927	Confirmation of wave-particle duality in electron diffraction experiments by Davisson-Germer and Thomson-Reid.
1928	Dirac's relativistic quantum theory: Introduced the Dirac equation and predicted antimatter.
1930	Dirac's quantum field theory: Established the foundation of quantum electrodynamics (QED).

interpreted probabilistically. Dirac accomplished the next conceptual step. This analysis teaches us how hard conceptual discoveries are. As Heisenberg remarked: "As a rule, new concepts come up in arather unclear and undeveloped form." This sequence of missedopportunities reminds us of how, in the technology space, IBMmissed Microsoft (i.e., creating a software giant), Microsoftmissed Apple (i.e., Apple products), Apple missed Google missed Facebook, and all of them missed OpenAI. Allwere gigantic missed opportunities. I wonder what else liesahead that we are missing now!

The early history of quantum mechanics illustrates how messy the discovery process really is. The textbooks and courses oftengloss over this aspect, presenting the final equations as if theywere reached clearly, smoothly, and logically. This is rarely thecase. They are often discovered through clever guesswork. Eventhe most beautiful Einstein field equations of gravity were discovered in this manner.60 I am reminded of a remark by HenriPoincare:61 "Guessing before proving! Need I remind you that it is so that all important discoveries have been made?"

Our analysis also reveals that the key challenges were conceptual rather than mathematical. Planck's

revolutionary quantum hypothesis is mathematically trivial: $\epsilon = hv$. Einstein's Nobel-winning equation is so simple that a high school student can understand: hv = W + Ek. Even Heisenberg-Born's matrixformulation or Schrödinger's equation is not tricky mathematically. Mathematical sophistication first emerged through Dirac's relativistic quantum mechanics and later in quantum field theory. Furthermore, the mathematical tools were already available and ready to be applied once the conceptual difficultieswere resolved. For example, matrices, probability theory, andpartial differential equations the main tools of quantum mechanics have been around for a long time. Similarly, for the theory of relativity. The mathematics of special theory is just elementary high school algebra, but the conceptual breakthroughs about space and time were colossal. The general theoryrequired more sophisticated mathematics, to be sure, but it was readily available, thanks to Riemann. 60The only instance in the history of physics where themathematical framework was also lacking, along with the needfor a conceptual breakthrough, was the discovery of the theory of gravitation. In addition to the conceptual breakthrough ofuniversal gravitation, Newton also had to develop themathematical tool needed, namely, the calculus. However, thisis the only exception that I am aware of. This analysis suggests another valuable lesson for the presenttime. Like the 1900s clouds, I believe we have a large cloud now n the horizon: the lack of a theory for deep neural networks large language models. By theory, I mean fundamental organizing principles that can predict important system-wideproperties, such as the structure and behavior of LLMs, fromtoken-level properties.62,63 To be sure, significant progress hasbeen made in the last three decades in neural network training, including the development of the backpropagation algorithm, various regularization techniques, reinforcement learning, andtransformer architecture, among others. However, these aremerely recipes for training; they do not provide a comprehensivetheory of deep neural networks or large language models(LLMs). This is the central conceptual challenge facing AI today.

In 1972, physics Nobel laureate Philip Anderson published an influential paper entitled "More is Different".64 He observed: The behavior of large and complex aggregates of elementary particles, it turns out, is not to be understood in terms of a simple extrapolation of the properties of a fewparticles. Instead, at each level of complexity, entirely newproperties appear, and the understanding of the newbehaviors requires research that we think is as fundamentalin its nature as any other ... At each stage, entirely new laws, concepts, and generalizations are necessary, requiring inspiration and creativity to just as great a degree as in the previous one. Psychology is not applied biology, nor biology applied chemistry."

In this sense, invoking another physics analogy, Newtonianmechanics and F = ma can explain the dynamics of a fewparticles. However, when we have Avogadro's number (6.02×1023) of molecules dynamically interacting in a gas, the collectivebehavior cannot be explained by applying Newton's law 1023times! To be sure, F = ma is going on at the molecular level, butmuch more happens at the system level that cannot beunderstood by Newton's Second Law alone. To explain macroscopic phenomena, we need entirely newconcepts, such as temperature, free energy, entropy, and chemical potential, to predict and explain the behavior of agas. These concepts are absent at the individual particle level in Newtonian mechanics. We require an entirely new conceptual and mathematical framework, known as statistical mechanics, toaddress this new physics. It turns out that we need the SecondLaw of Thermodynamics and not the Second Law of Newton. This dichotomy between classical and statistical mechanics islike the proverbial "seeing trees but not the forest". The F = maperspective is "seeing the trees," and $S = k \ln W$ is "seeing theforest." Likewise, large language models are not mere stochasticautocomplete engines. They have new emergent capabilities that require creating a new conceptual framework similar to the transformation from Newtonian to statistical mechanics or from classical to quantum mechanics. The LLMs may not have developed a human-like understanding of their domain, but they seem to have

acquired a different kind of understanding and intelligence. Although it is difficult to say without anyuncertainty that AI is at a "1900-moment," the signs are compelling. For millennia, we have taken for granted themeanings of words such as "understanding" and "intelligence" without much introspection. With the advent of LLMs, we are compelled to reevaluate our understanding of such concepts. LLMs raise profound philosophical questions about consciousness, free will, and the nature of creativity and intelligence, conceptual questions with which we are only beginning tograpple.

So, what would a mathematical theory of LLMs look like? As noted, I believe mathematical tools are already available: linearalgebra, probability theory, statistical mechanics, game theory, graph theory, group theory, and topology. The challenge lies indiscovering new concepts necessary for this problem. Asdiscussed, quantum theory was born from the analysis of theenergy distribution in blackbody radiation. Classical physicsbased theories could not explain this distribution, which compelled Planck to propose a quantum hypothesis. Similarly,in well-trained deep neural networks, the connection weights are distributed lognormally. Neither the Hopfield nor the Boltzmann Machine model, which were recognized with the 2024 Nobel Prize in Physics, can predict or explain the lognormal outcome. Recently, a new conceptual framework, 63 called statistical teleodynamics, which combines game theory and statistical mechanics, has been proposed to predict this outcomeas a first step toward a mathematical theory of LLMs. Borrowing from physics, the Hopfield and Boltzmann machine models employ energy minimization, where as the new framework uses effective-utility maximization from economics as its organizing principle.

The ultimate theory of LLMs can potentially upend our views of cognition and sentience, much like the "1900-moment" did inphysics a century ago. Thus, as Planck and Heisenberg remarkedabout how new concepts are born amid profound confusion, understanding the historical evolution of the quantummechanical concepts could be helpful in a similar situation tothat we face in artificial intelligence.

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